

Laser and Accelerator Tech for High-Brightness Laser-Compton Light Sources

**KEK Satellite Meeting
Tokai-mura, Japan**



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January 31, 2014**

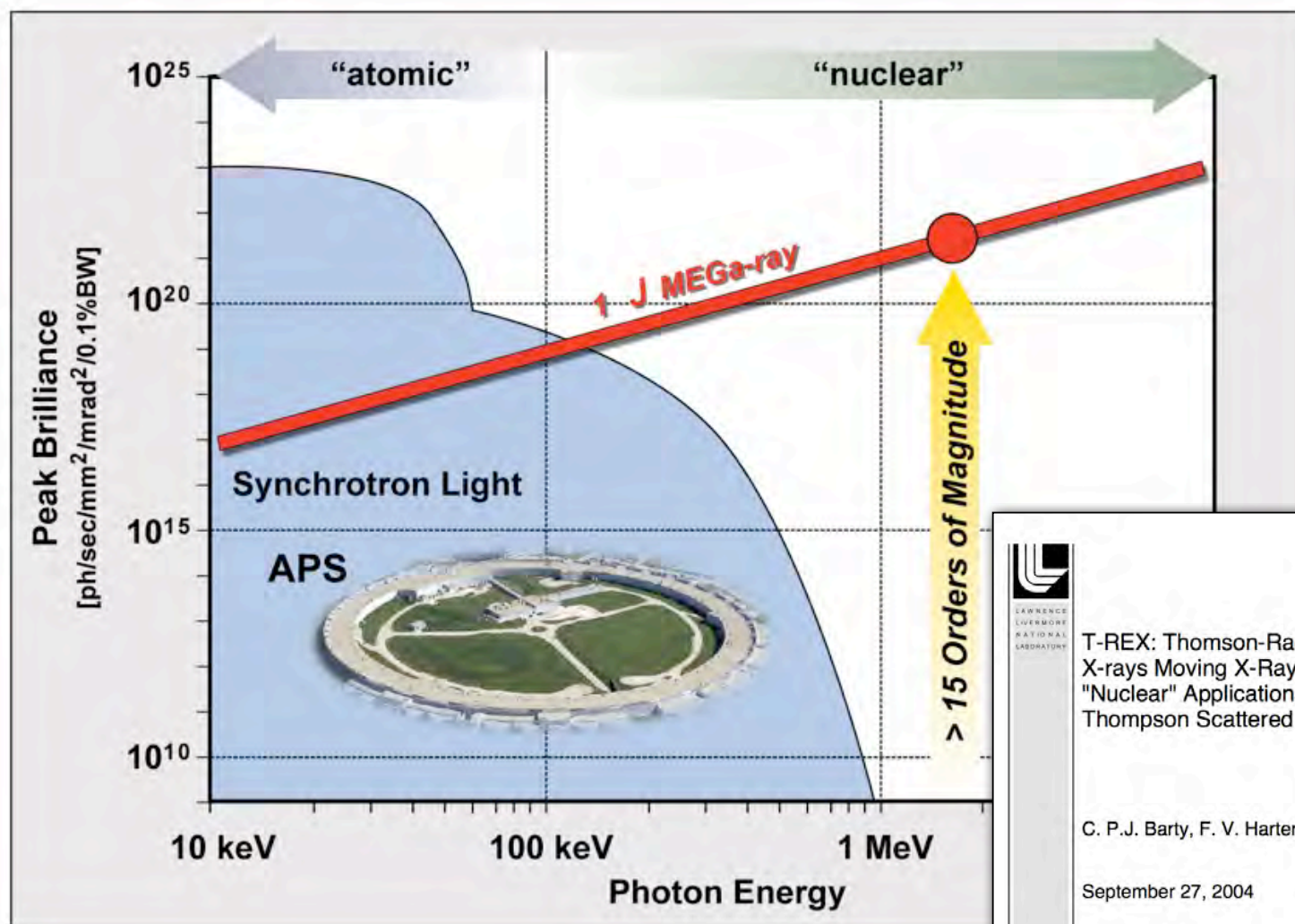
The MEGa-ray and Nuclear Photonics efforts described in this presentation represent contributions from 11 institutions

Marvin	Adams	TAMU
Chris	Adolphsen	SLAC
Felicie	Albert	LLNL
Gerry	Anderson	LLNL
Scott	Anderson	LLNL
Paul	Armstrong	LLNL
Chris	Barty	LLNL
Andy	Bayramian	LLNL
Bret	Beck	LLNL
Glenn	Beer	LLNL
Shawn	Betts	LLNL
Dave	Boyle	TAMU
Patrick	Brantley	LLNL
Eugene	Brooks	LLNL
Arno	Candel	SLAC
Bill	Charlton	TAMU
Sam	Chu	SLAC
Eric	Cormier	UBordeaux
Rick	Cross	LLNL
Dan	Cutiou	ELI-NP
Gary	Deis	LLNL
Bob	Demaret	LLNL

Shawn	Densberger	LLNL
Valery	Dolgashev	SLAC
Chris	Ebbers	LLNL
Mike	Fazio	SLAC
Diana	George	LLNL
David	Gibson	LLNL
Marc	Gunther	LMU
Dietrich	Habs	LMU
Chris	Hagmann	LLNL
Ryoichi	Hajima	JAEA
James	Hall	LLNL
Fred	Hartemann	LLNL
Corrine	Izak	CEA
Michael	Jentschel	ILL
Micah	Johnson	LLNL
Ed	Jones	LLNL
Erik	Jongewaard	SLAC
Zenghai	Li	SLAC
Cecile	Limborg-Deprey	SLAC
Roark	Marsh	LLNL
Scott	McKinley	LLNL
Dennis	McNabb	LLNL
Jim	Morel	TAMU

Ed	Morse	UCB
Kaila	O'Neil	LLNL
Henry	Phan	LLNL
Norbert	Pietralla	GSI
John	Post	LLNL
Matt	Prantil	LLNL
Cesar	Pruneda	LLNL
Sofia	Quagllioni	LLNL
Tor	Raubenheimer	SLAC
Vladimir	Semenov	LLNL
Michio	Seya	JAEA
Rich	Shuttlesworth	LLNL
David	Stevens	LLNL
Sami	Tantawi	SLAC
Peter	Thiorlf	LMU
Arnold	Vlieks	SLAC
Faya	Wang	SLAC
Juwen	Wang	SLAC
Caroline	Winters	LLNL
Sheldon	Wu	LLNL
Victor	Zamfir	ELI-NP
Feng	Zhou	SLAC

The characteristics of optimized laser-Compton gamma-ray sources enable “nuclear photonics”



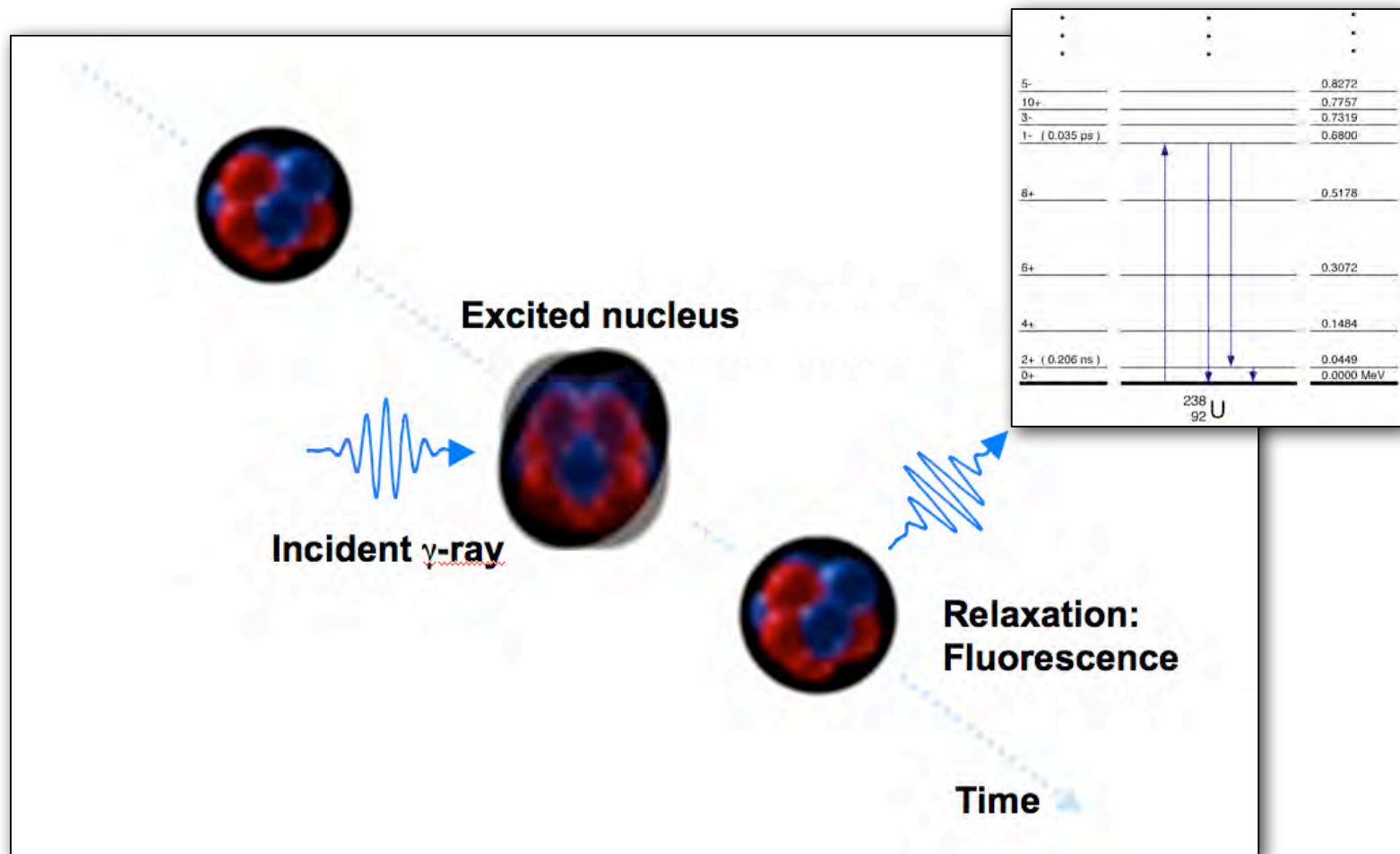
UCRL-TR-206825

T-REX: Thomson-Radiated Extreme X-rays Moving X-Ray Science into the "Nuclear" Applications Space with Thompson Scattered Photons

C. P.J. Barty, F. V. Hartemann

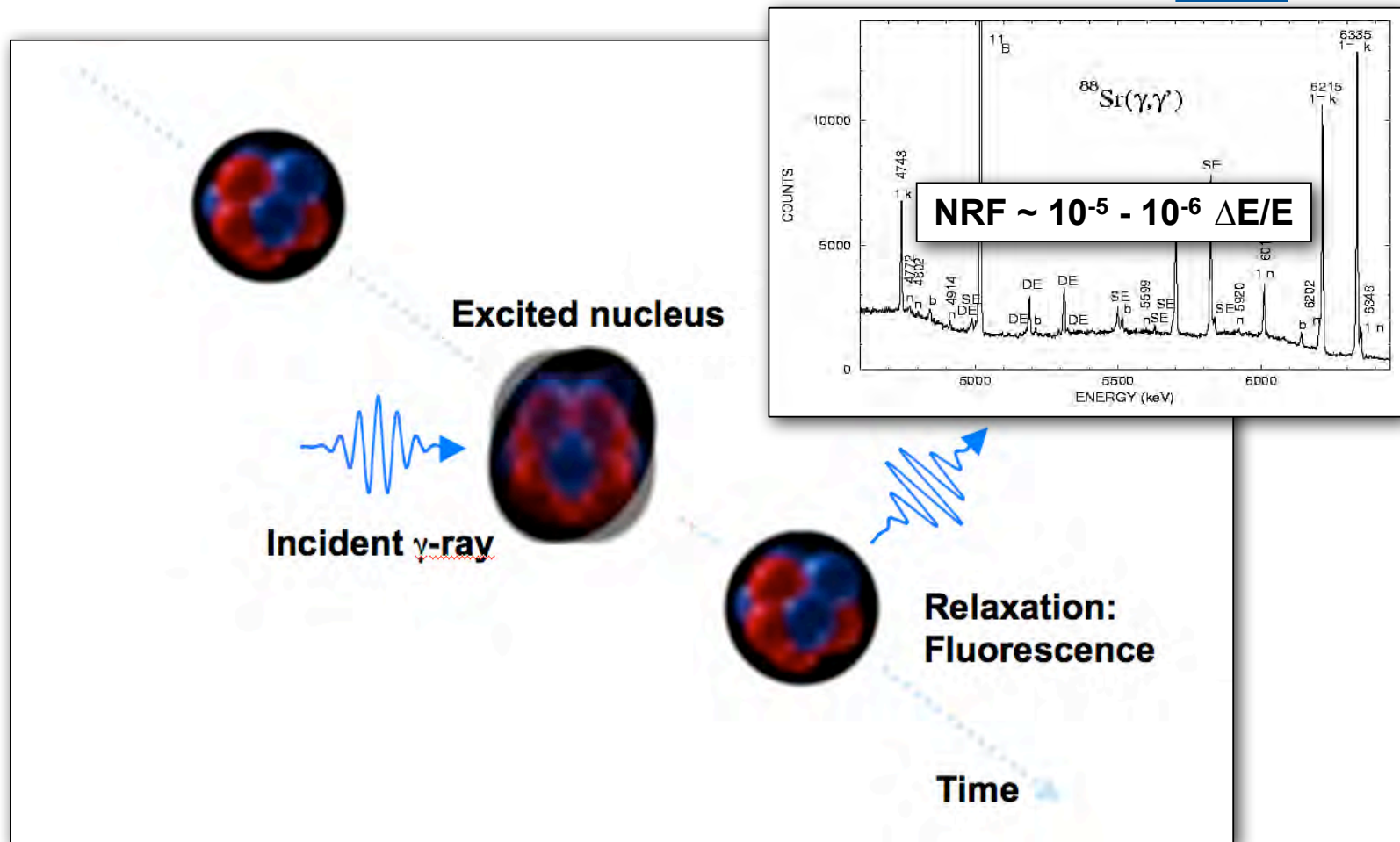
September 27, 2004

Nuclear resonance fluorescence is easily excited narrowband laser-Compton sources



Nuclear Resonance Fluorescence depends upon the number of protons and the number of neutrons in the nucleus and is an isotope-specific material signature

Intrinsic NRF widths are of order meV but are thermally (Doppler) broadened to of order eV



Selective excitation of NRF transitions is possible with laser-Compton gamma-ray source bandwidths of order $\Delta E/E \sim 10^{-3}$

Potential NRF-based Applications of Bright Gamma Sources are Numerous



HEU Grand Challenge
detection of shielded material



Nuclear Fuel Assay
100 parts per million per isotope



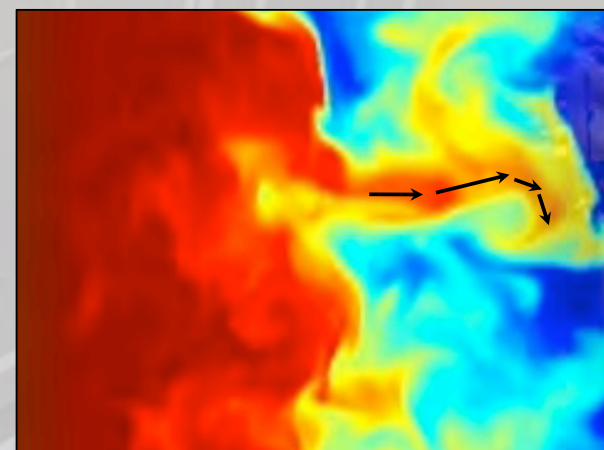
Waste Imaging & Assay
non-invasive content certification



Industrial NDE
micron-scale & isotope specific

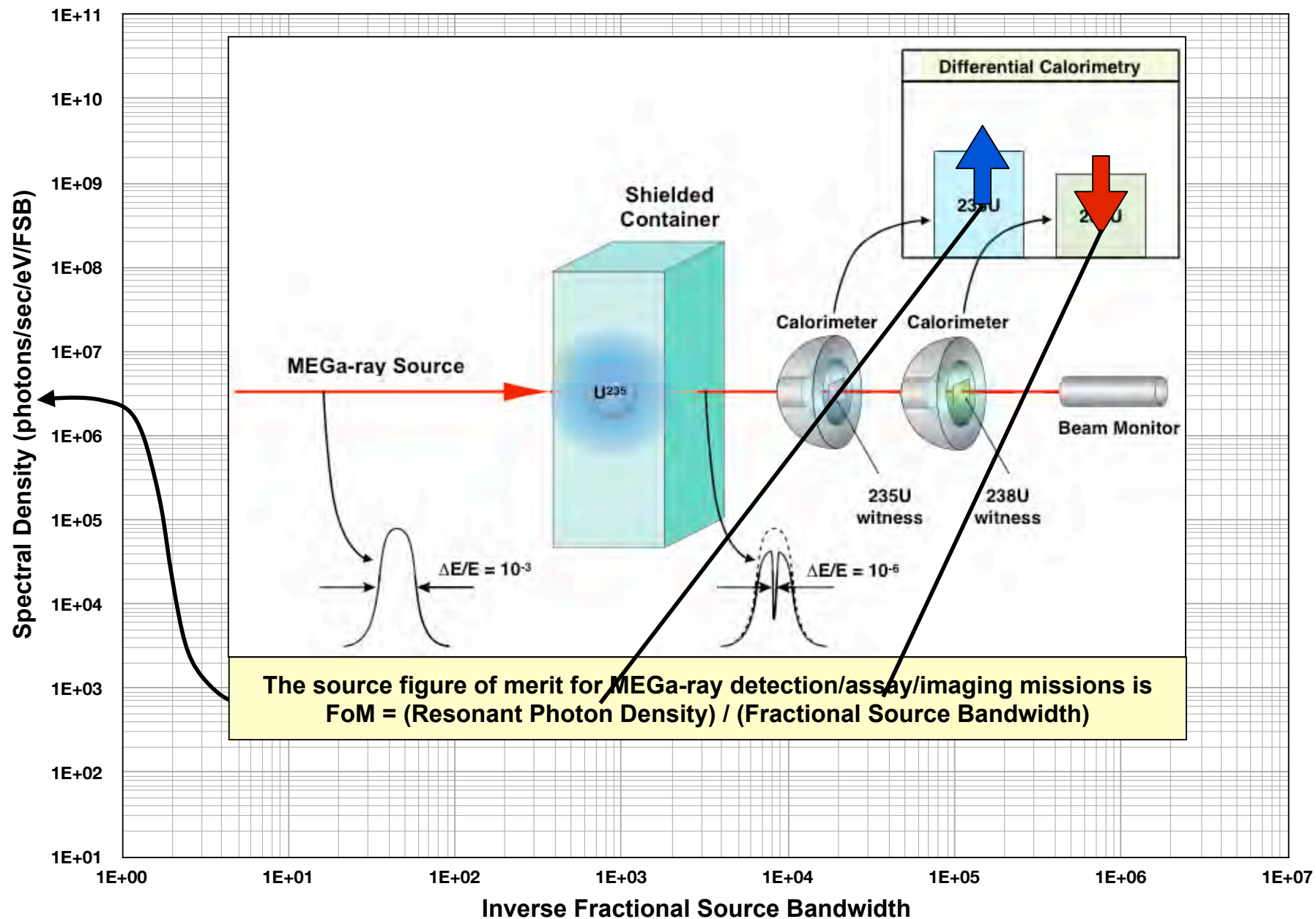


Medical Imaging
low density & isotope specific

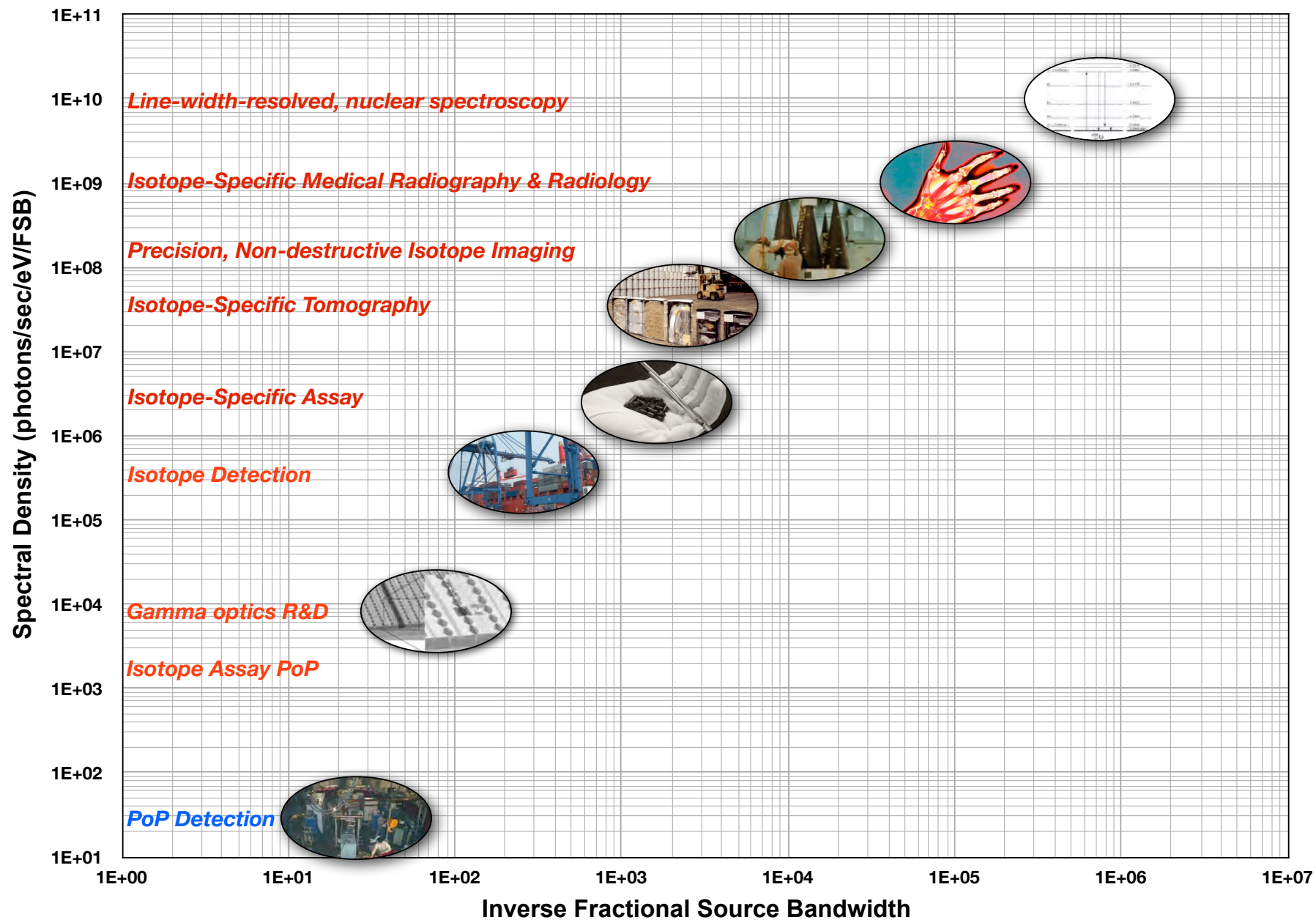


Dense Plasma Science
isotope mass, position & velocity

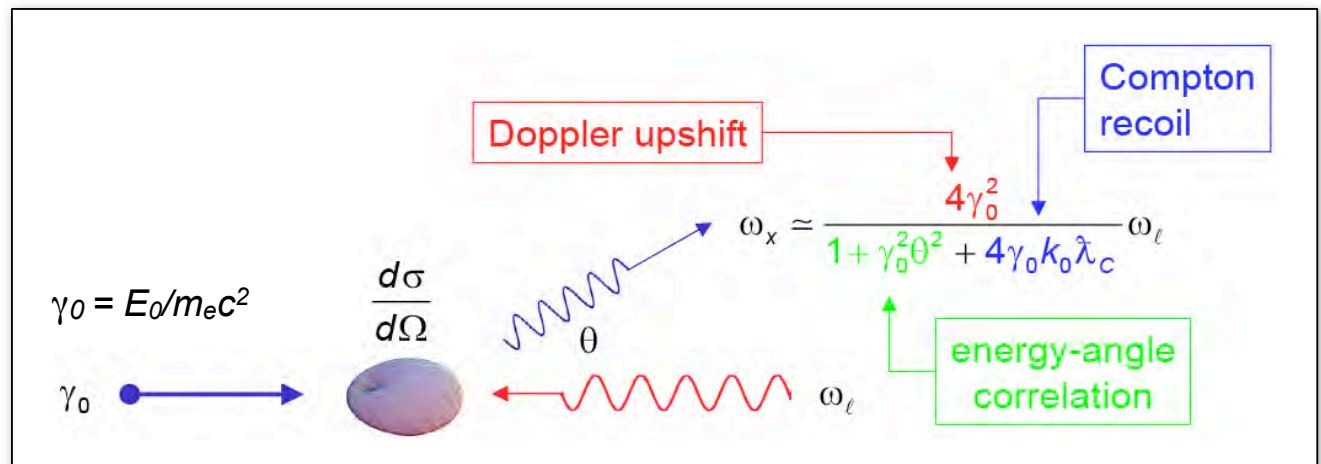
Spectral Density is the key source metric for most NRF applications



New applications become viable with increasing Spectral Density



Laser Compton back scattering off of high energy electrons can produce tunable x-ray & gamma-rays



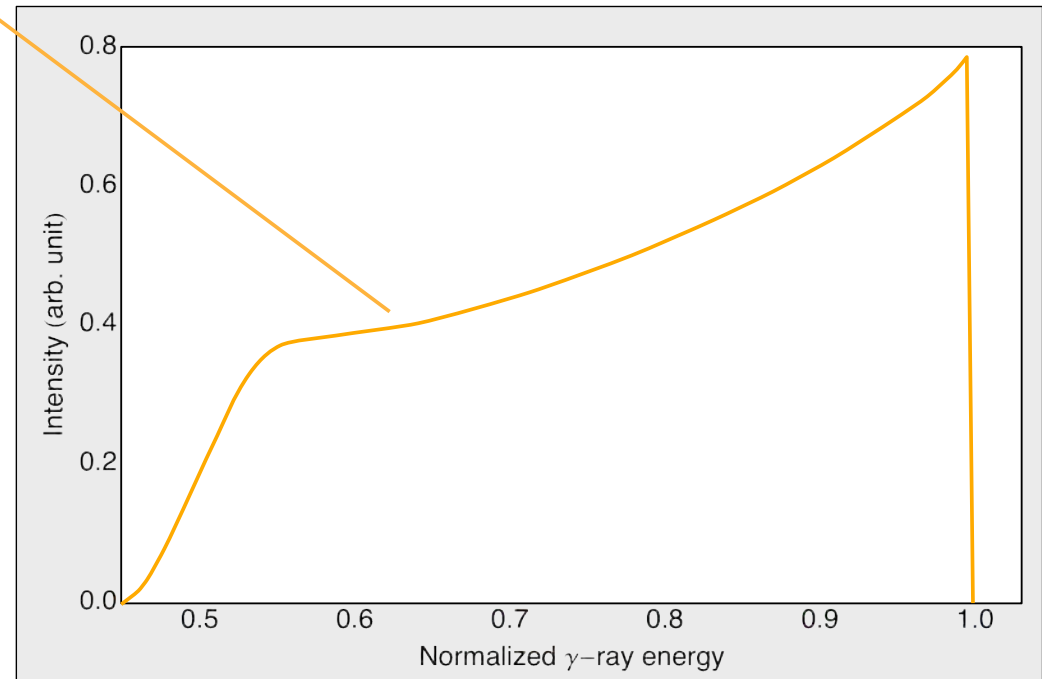
A photon flux of 1.7×10^{24} ph/cm² = “unity” efficiency

1.7×10^{24} ph/cm² @ 532 nm in a 100 micron spot = 44 J!

Overall Compton scattering is broadband, but it is highly angle correlated and is 'narrowband' on axis



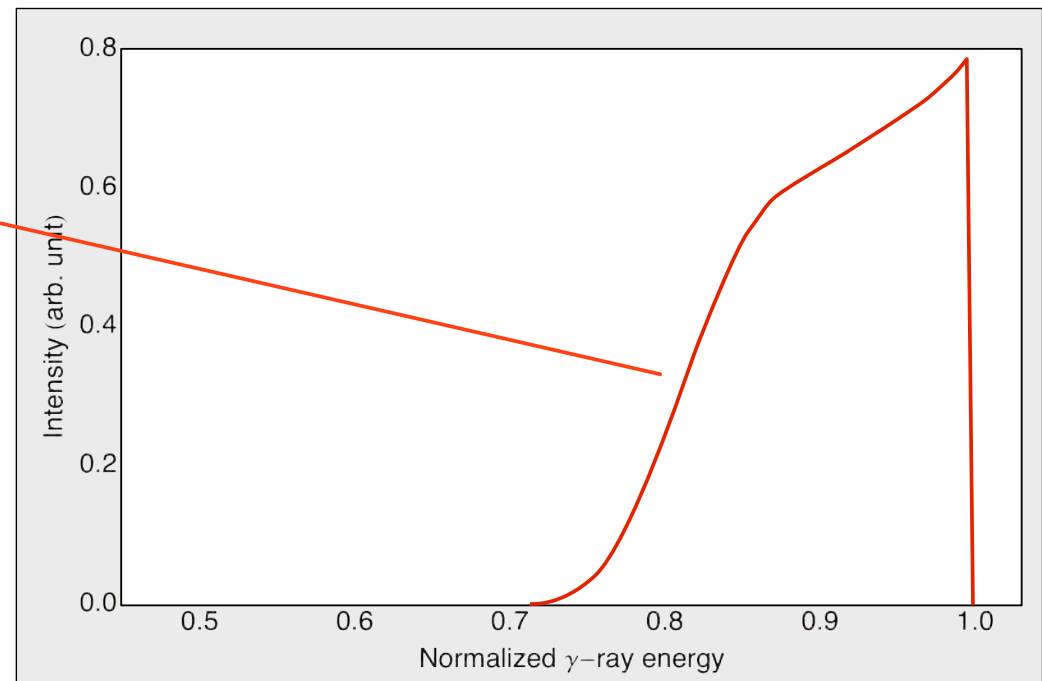
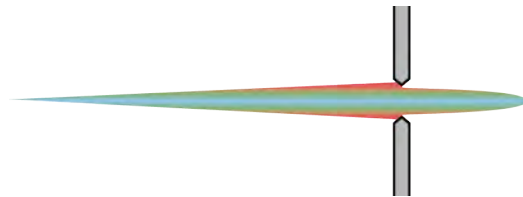
$\Delta\Omega ; \pi \left(\frac{1}{\gamma} \right)^2$
few mrad



Overall Compton scattering is broadband, but it is highly angle correlated and is 'narrowband' on axis



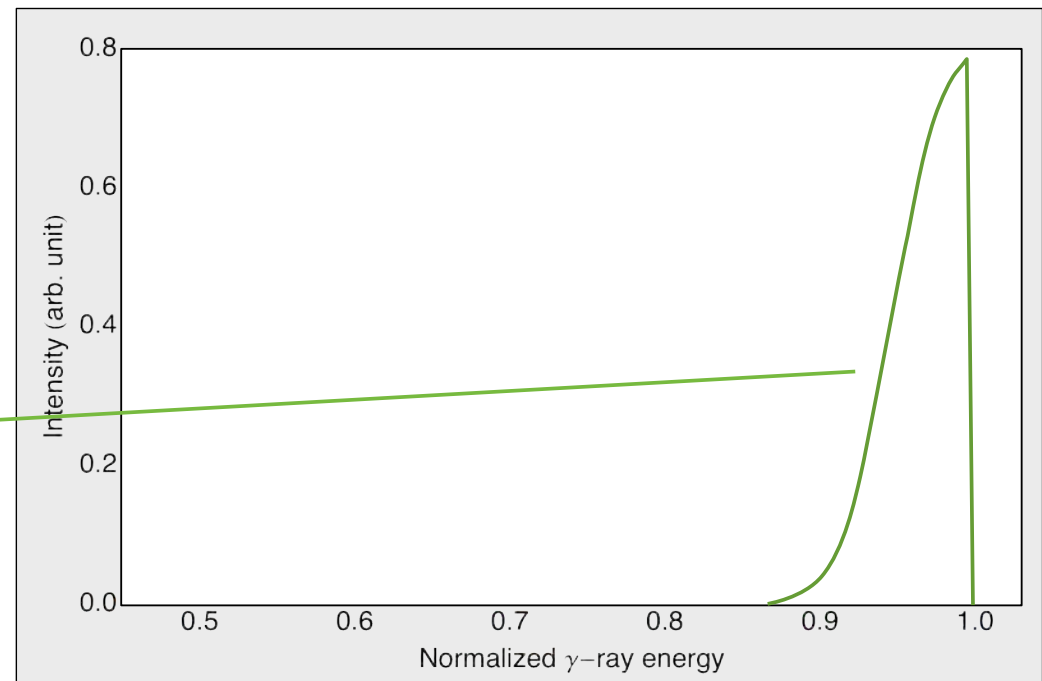
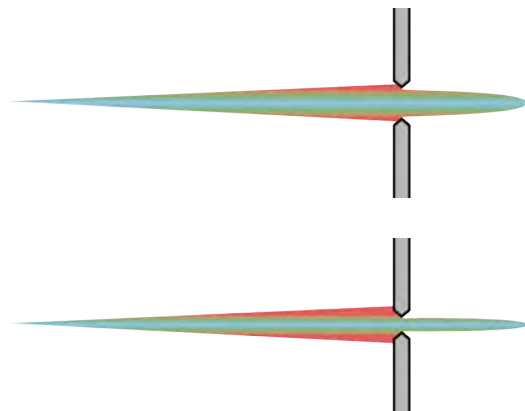
$\Delta\Omega ; \pi \left(\frac{1}{\gamma} \right)^2$
few mrad



Overall Compton scattering is broadband, but it is highly angle correlated and is 'narrowband' on axis



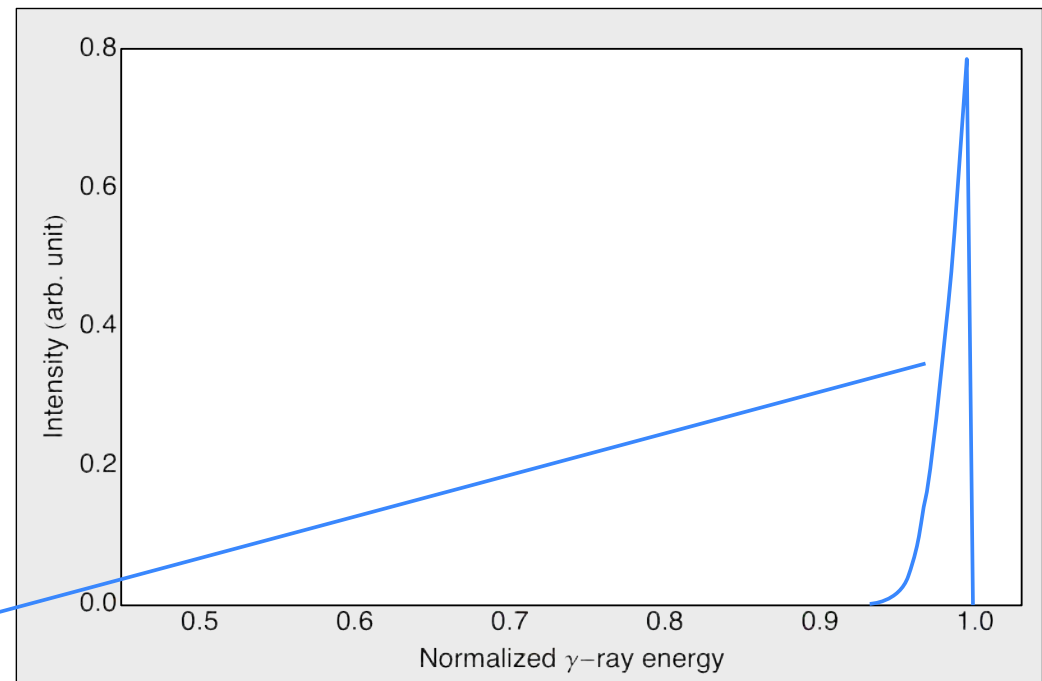
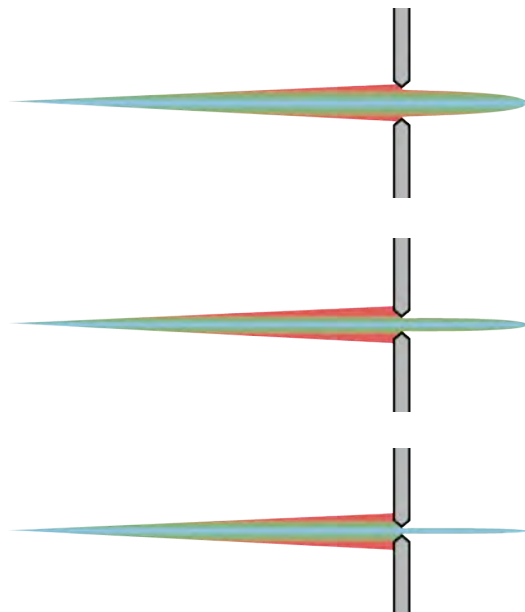
$\Delta\Omega ; \pi \left(\frac{1}{\gamma} \right)^2$
few mrad



Overall Compton scattering is broadband, but it is highly angle correlated and is 'narrowband' on axis



$\Delta\Omega ; \pi \left(\frac{1}{\gamma} \right)^2$
few mrad

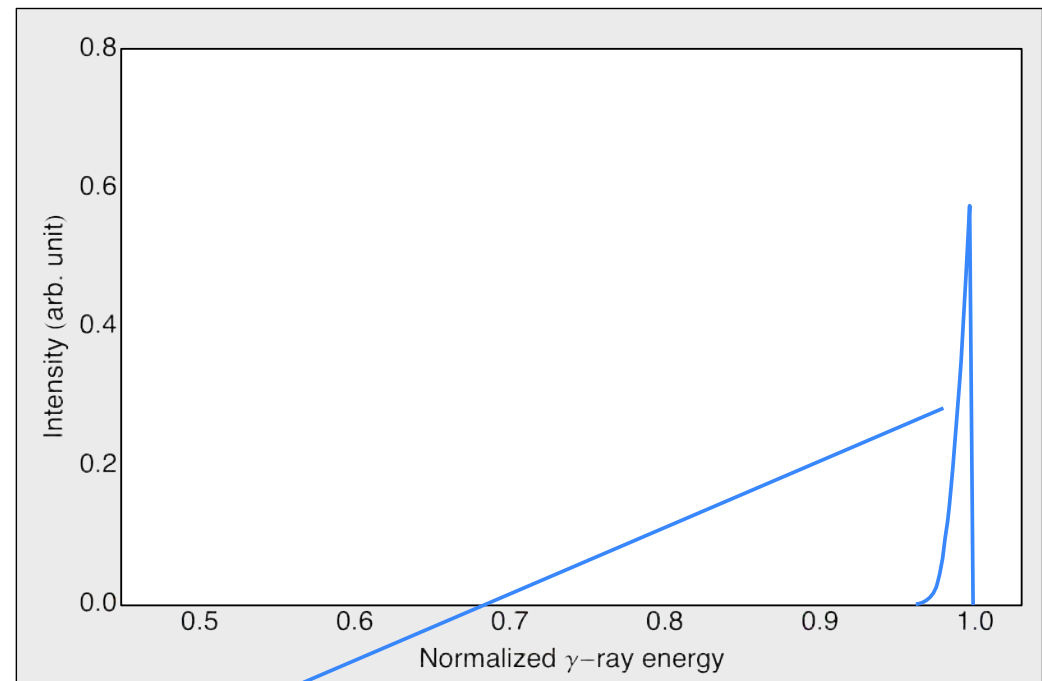
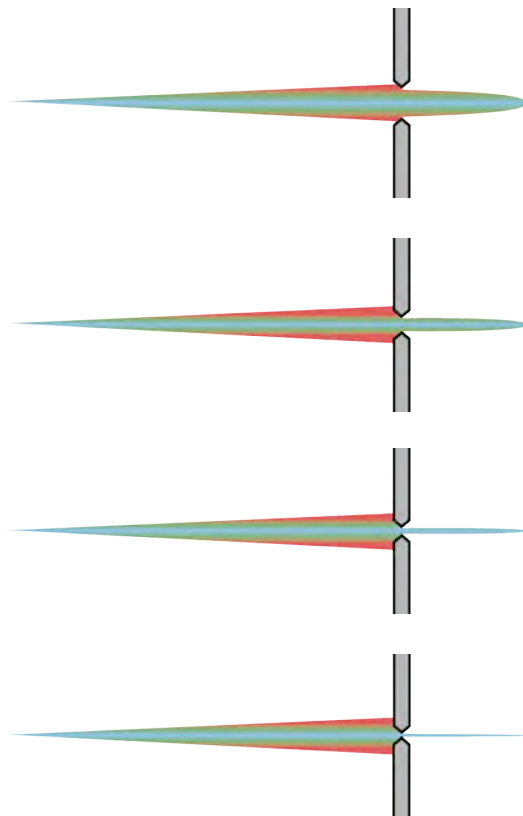


Overall Compton scattering is broadband, but it is highly angle correlated and is 'narrowband' on axis



$$\Delta\Omega ; \pi \left(\frac{1}{\gamma} \right)^2$$

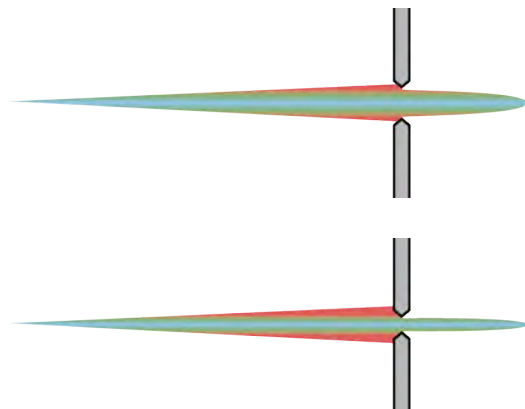
few mrad



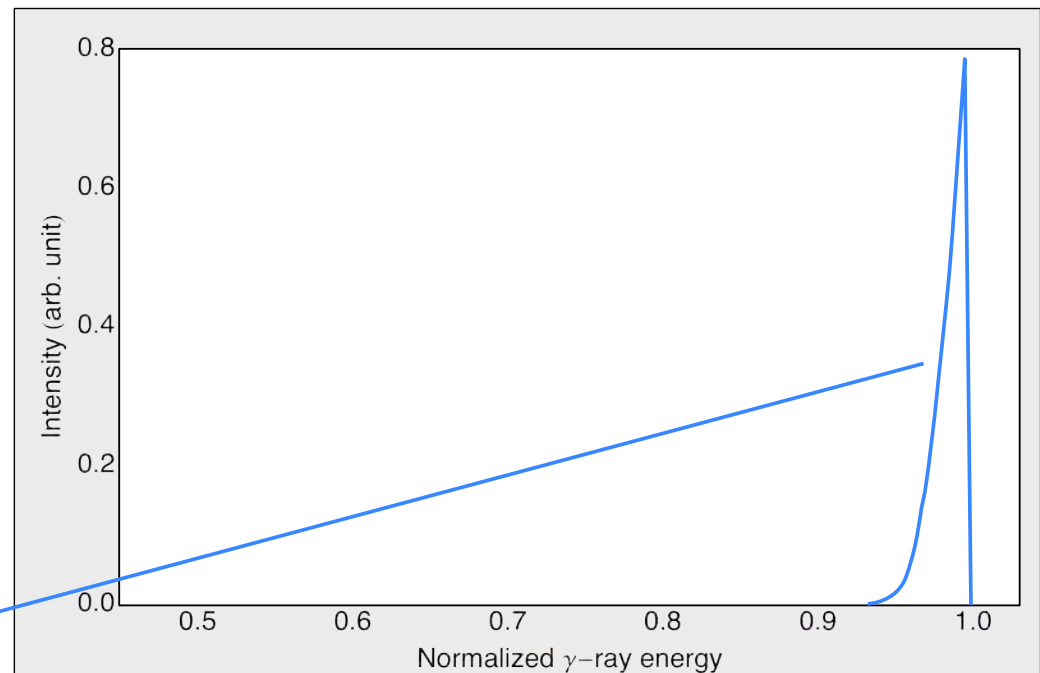
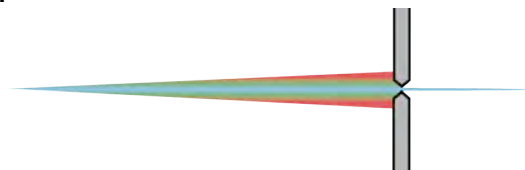
Overall Compton scattering is broadband, but it is highly angle correlated and is 'narrowband' on axis



$\Delta\Omega ; \pi \left(\frac{1}{\gamma} \right)^2$
few mrad



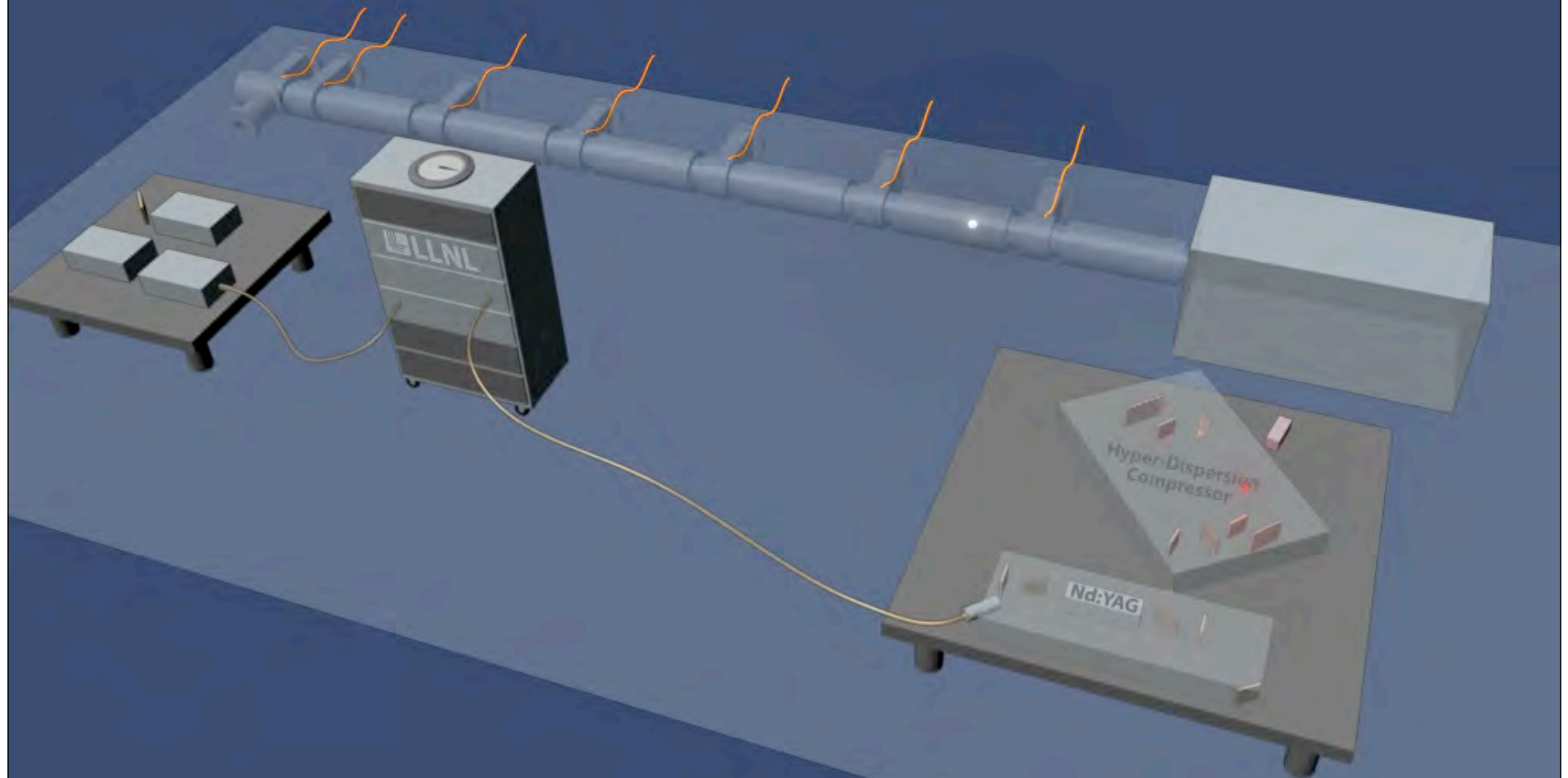
$\Delta\Omega ; \pi \left(\frac{\varepsilon_n}{\gamma\sigma} \right)^2$
few μ rad



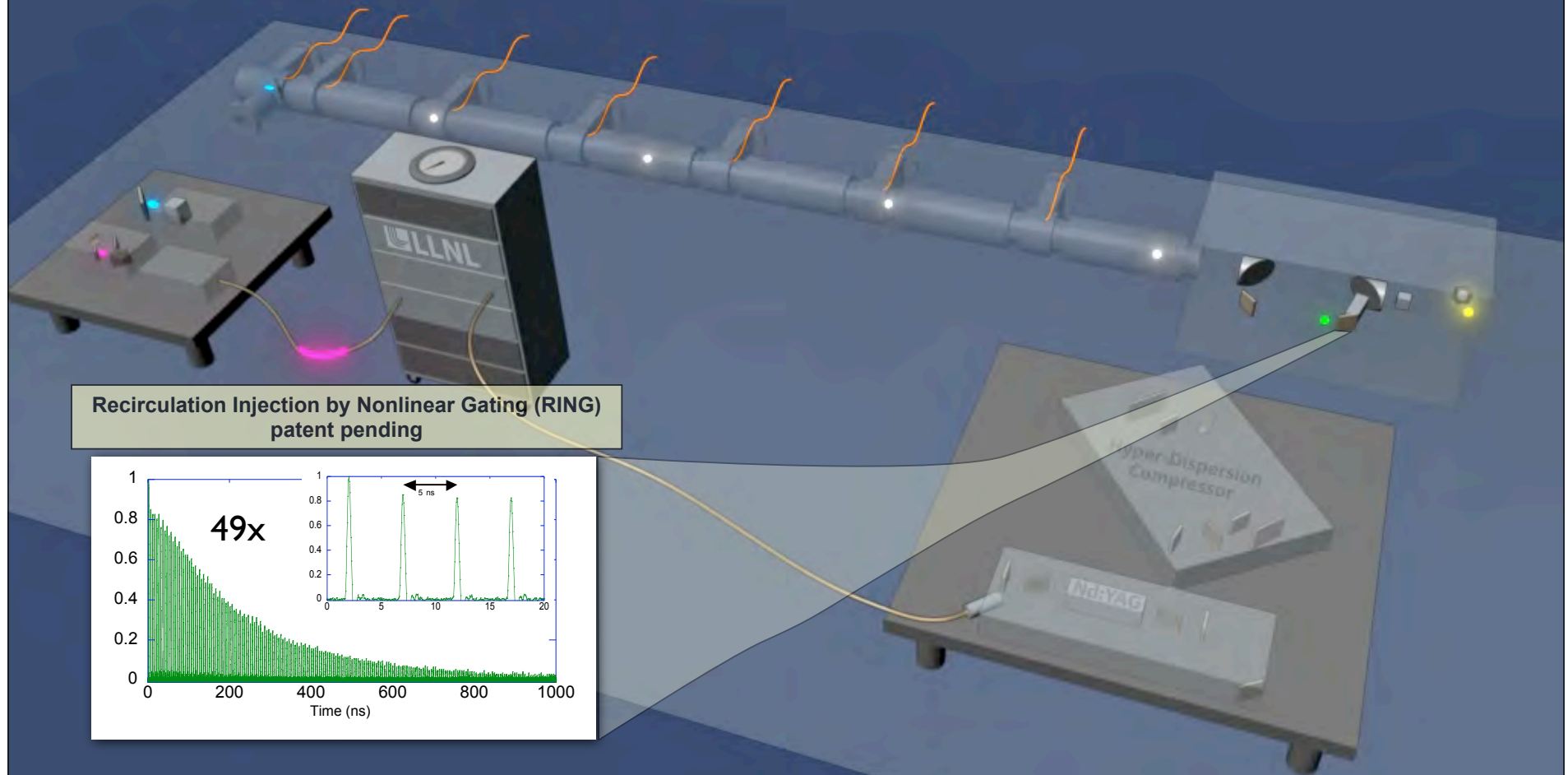
“Mono-Energetic Gamma-rays” - MEGa-rays

Optimally designed sources can have fractional bandwidths of $\sim 10^{-3}$ FWHM

High-flux, laser-Compton scattering arrangements aim to produce high photon & electron densities at a common focus

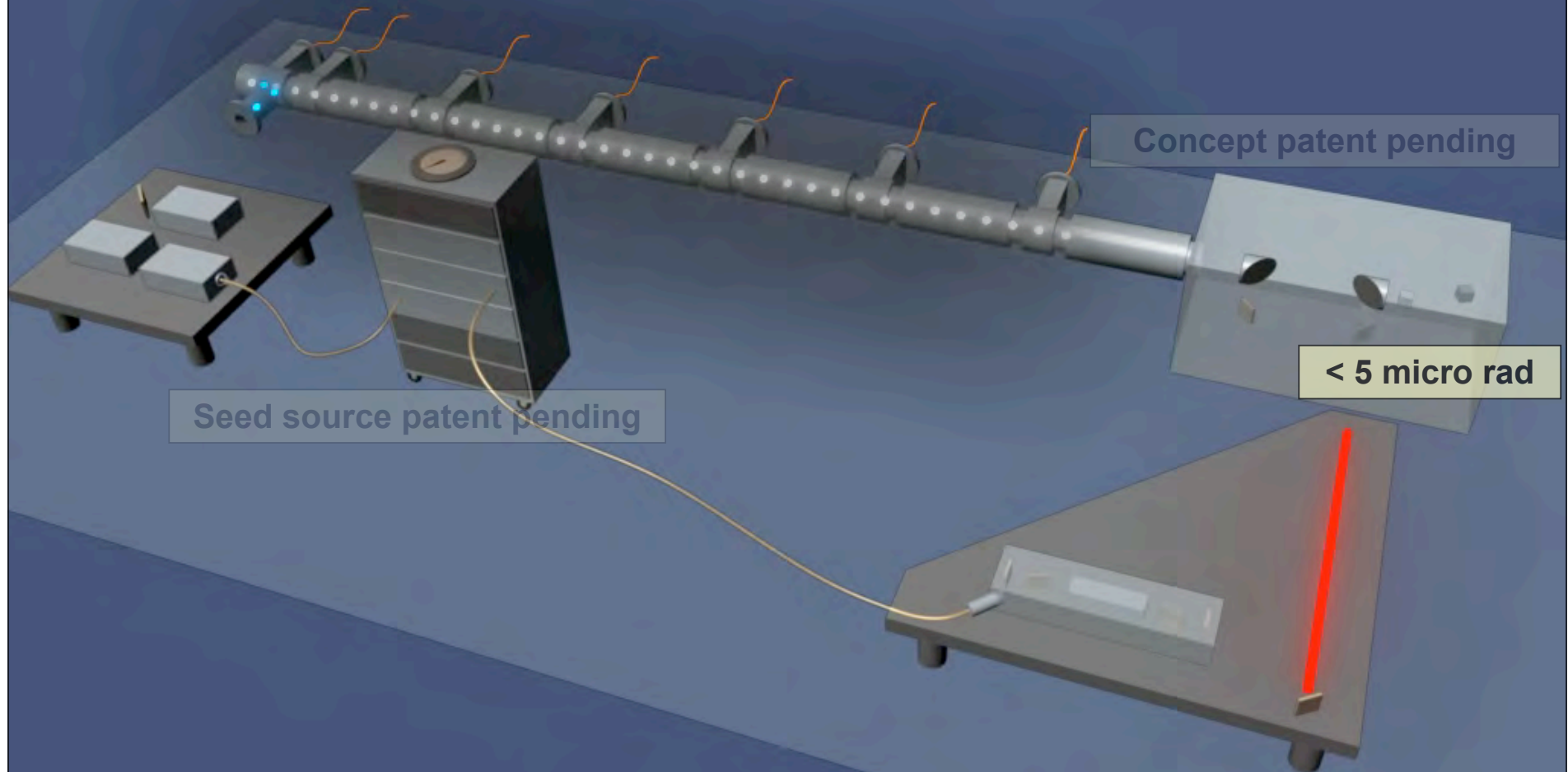


High-flux, laser-Compton scattering arrangements perturb the laser pulse energy very little during the interaction



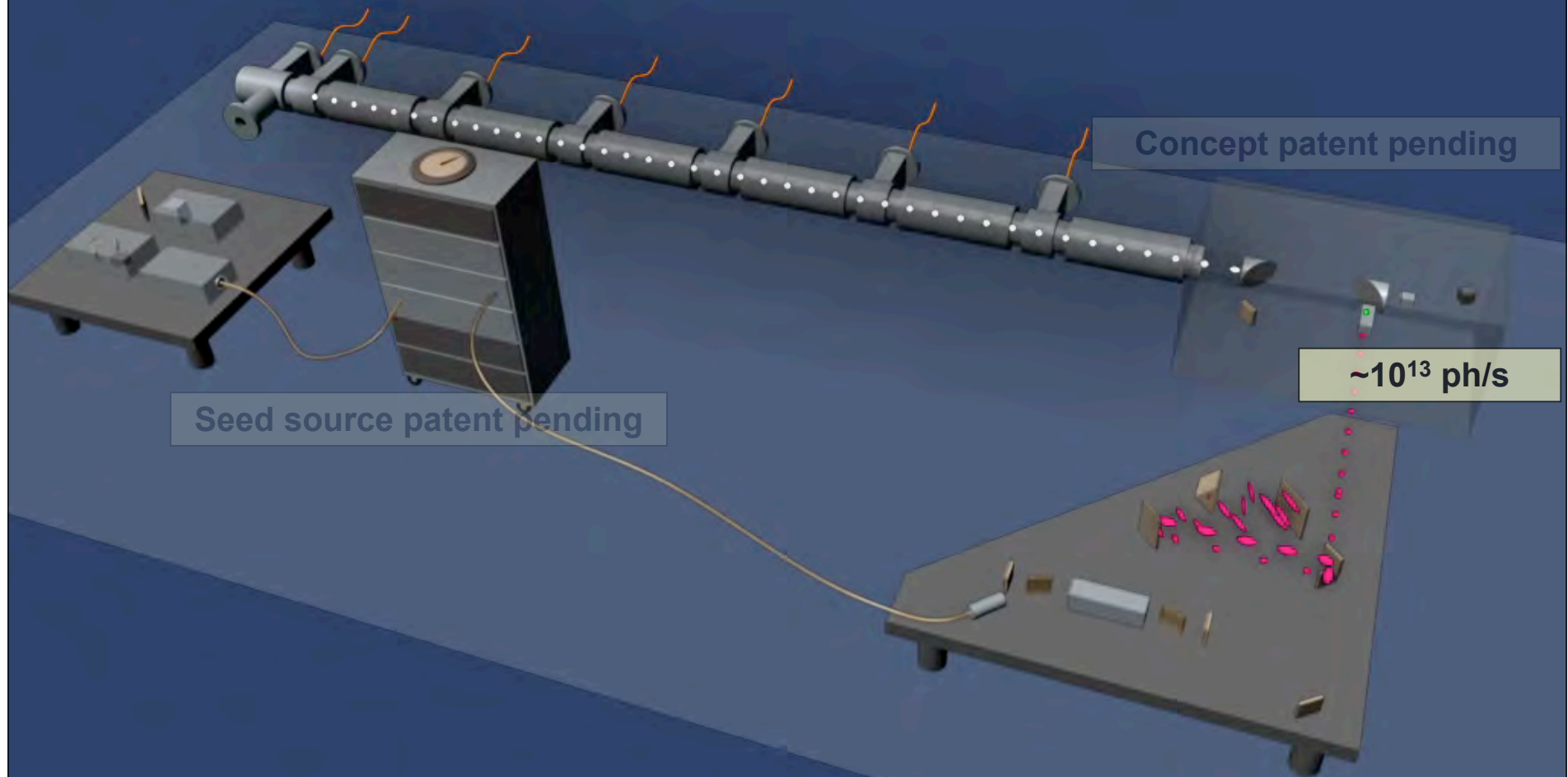
**Recirculation can give > 20x increase in Compton photon production for “free”
RING positioning requirements are 10,000x less stringent than Cavity schemes**

2011 multi-GHz, multi-bunch laser-Compton source concept



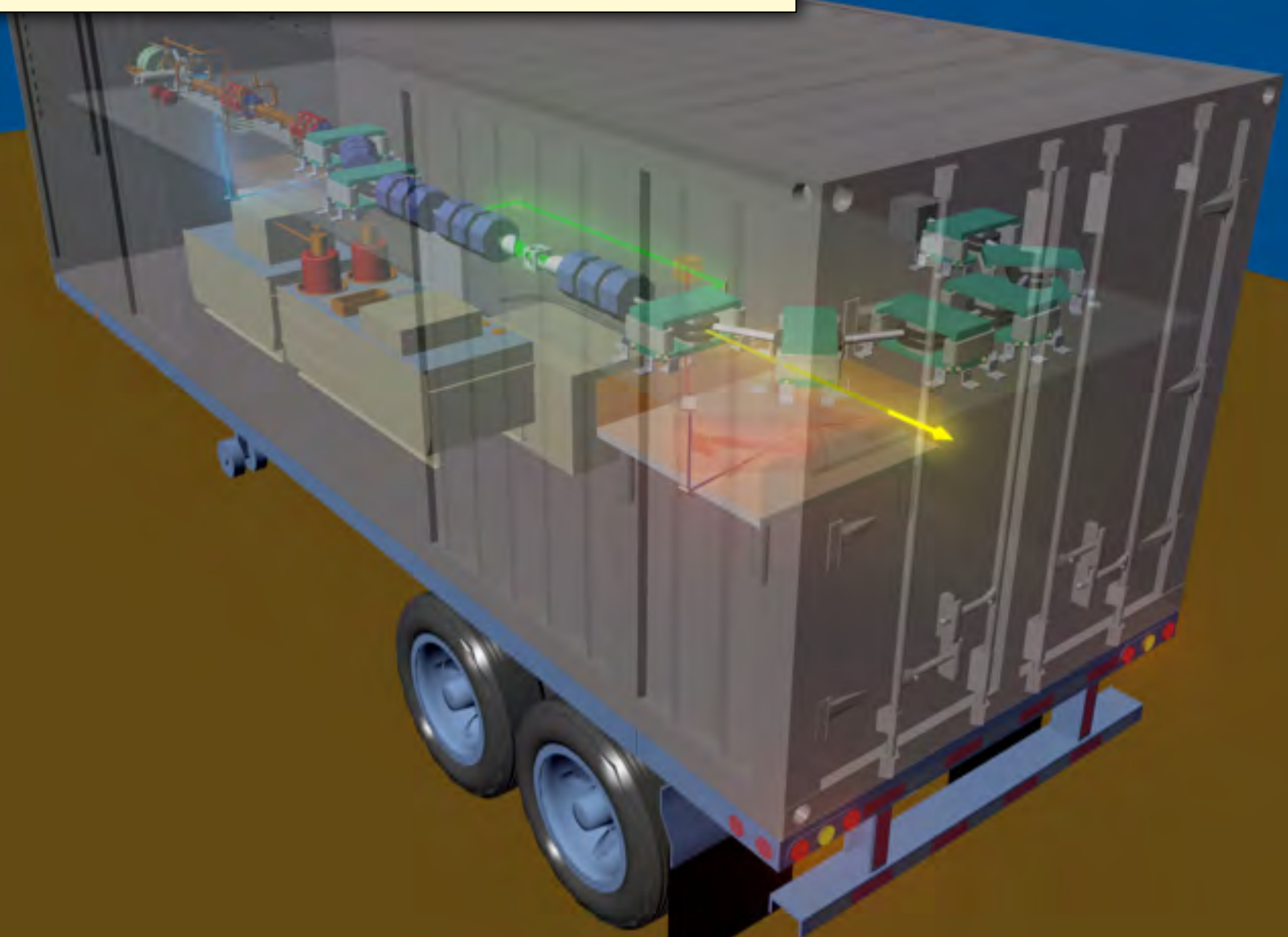
Highly collimated - reduces bandwidth, complexity of photon gun drive laser, interaction laser and system timing but requires high energy laser

LLNL's "Picket Fence" multi-GHz, laser-Compton source concept



This configuration enables near "unity" efficiency, operates with high beam current, minimizes bandwidth and is intrinsically synchronized to RF clock

Many applications also require that the underlying technologies be robust and compact





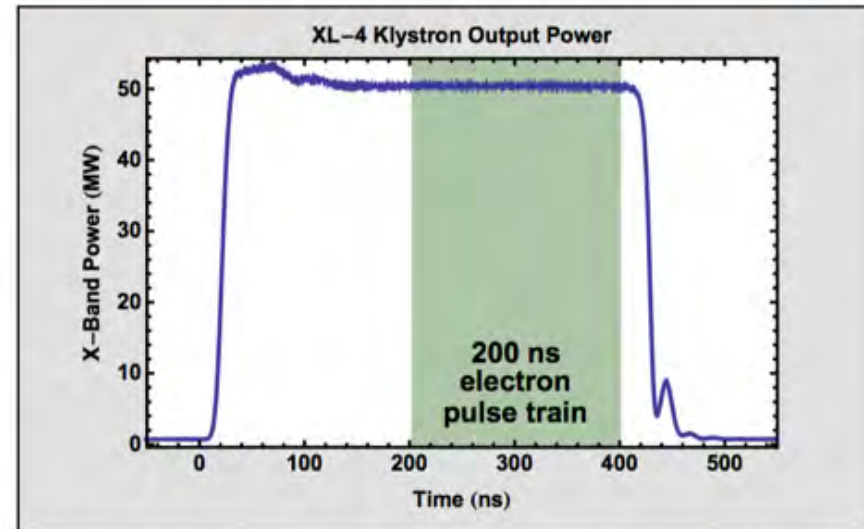
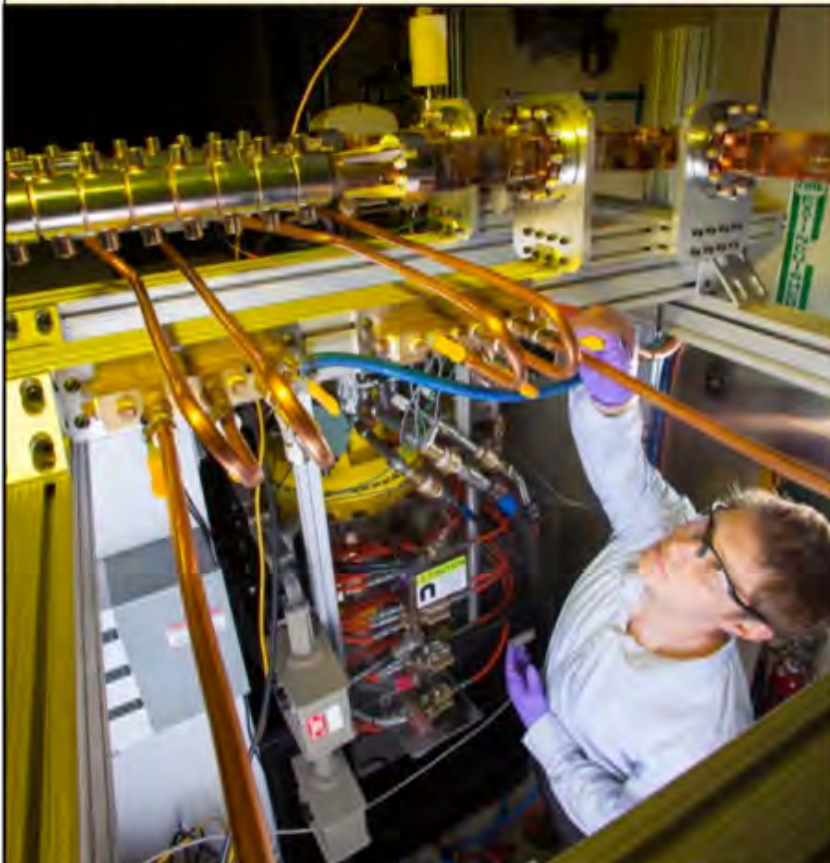
B194 X-Band Test Station

LLNL has designed & constructed a compact x-band accelerator in order to develop & demonstrate advanced, high-flux, laser-Compton x-ray & gamma-ray architectures

RF Power combines the best of SLAC klystron technology & commercial solid state modulators



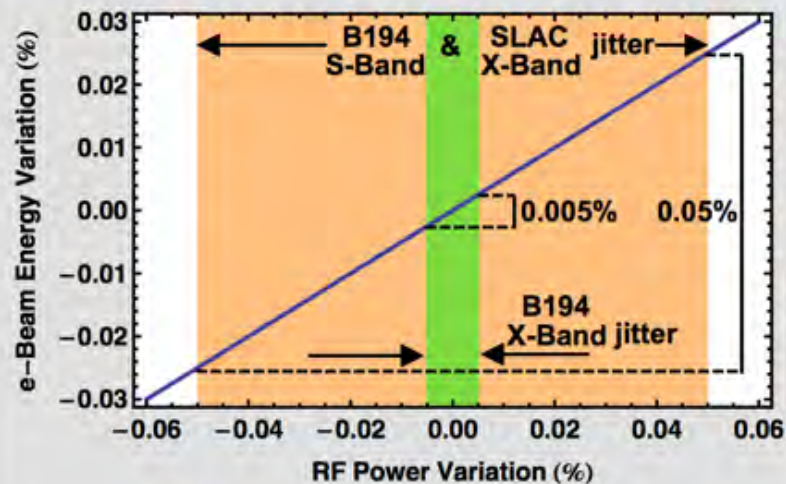
**State-of-the-art solid-state, 420 kV,
300 A Modulator and 50 MW Klystron**



Performance of the XL4 klystron and ScandiNova modulator exceed all of our requirements

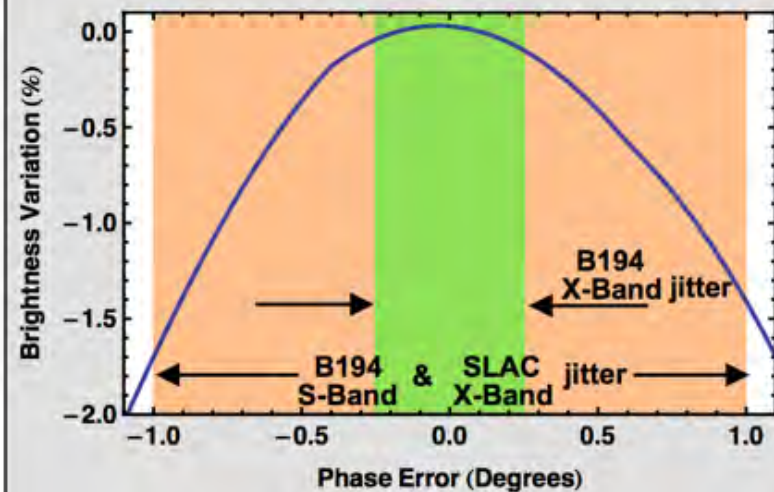


E-Beam Energy vs. RF Power



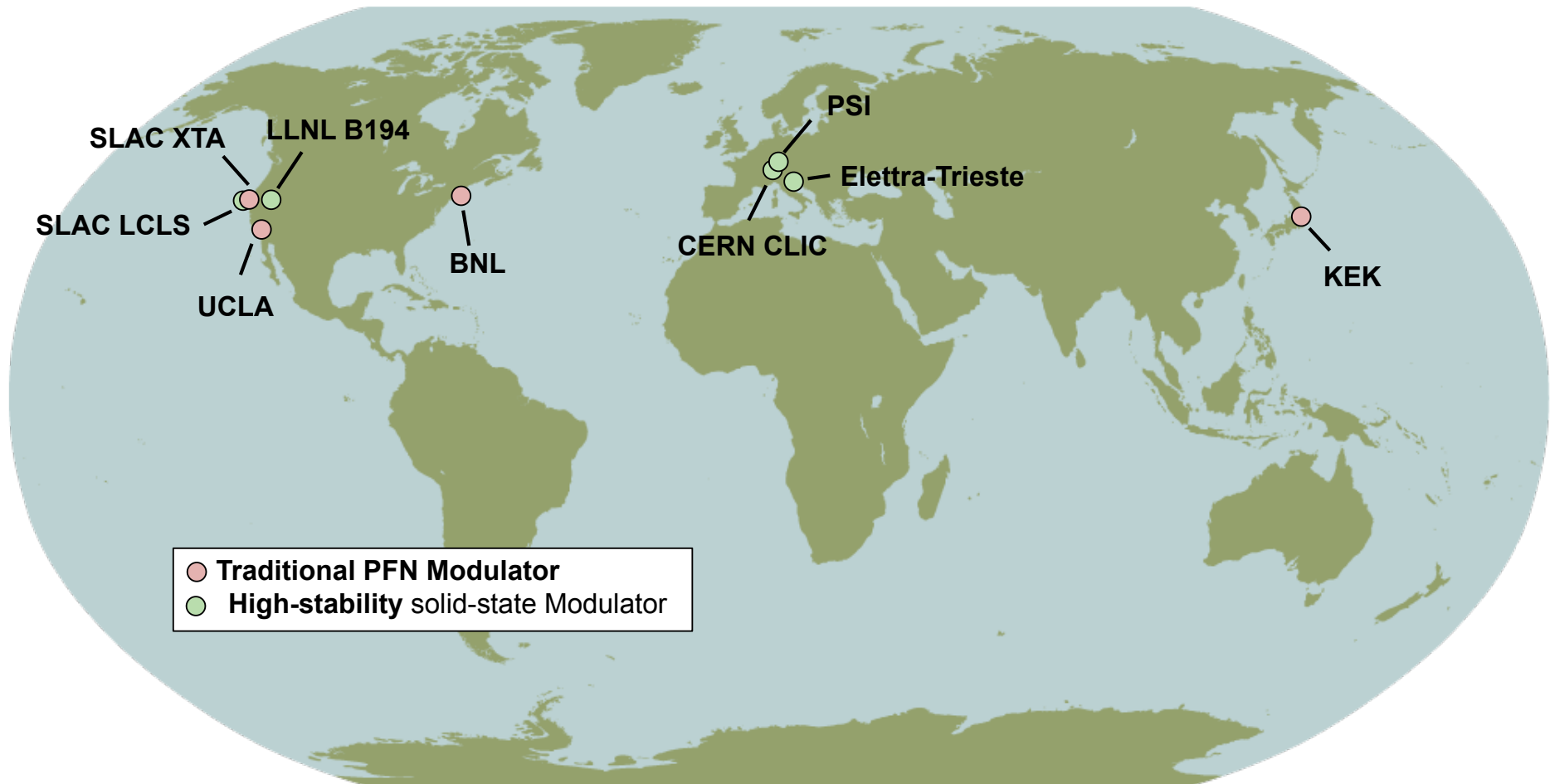
This system is 10x more stable in amplitude than existing SLAC and LLNL sources.

E-Beam Brightness vs. RF Phase

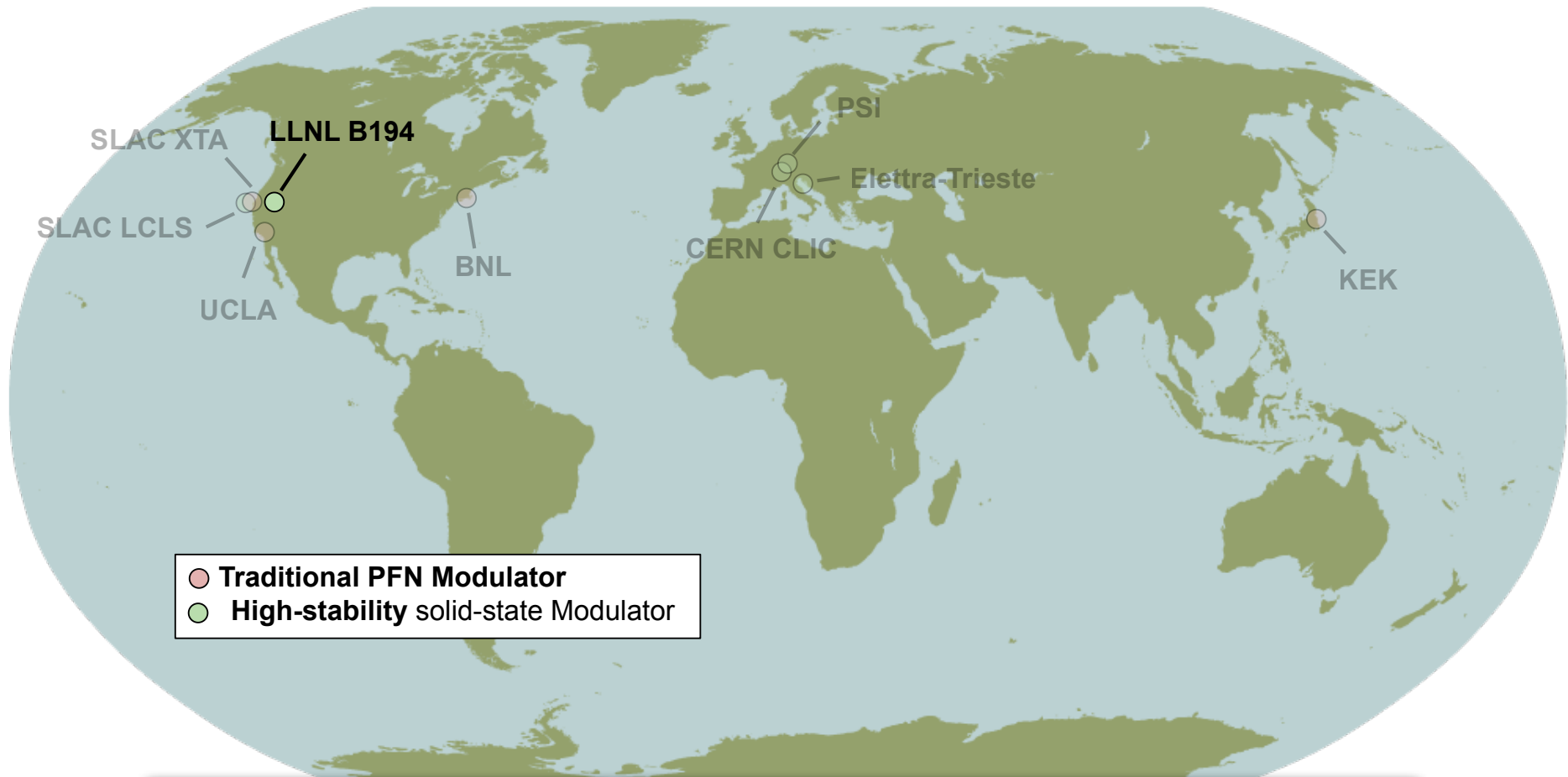


4x improvement in RF phase stability corresponds to 10x improvement in brightness stability.

Worldwide high power x-band sources

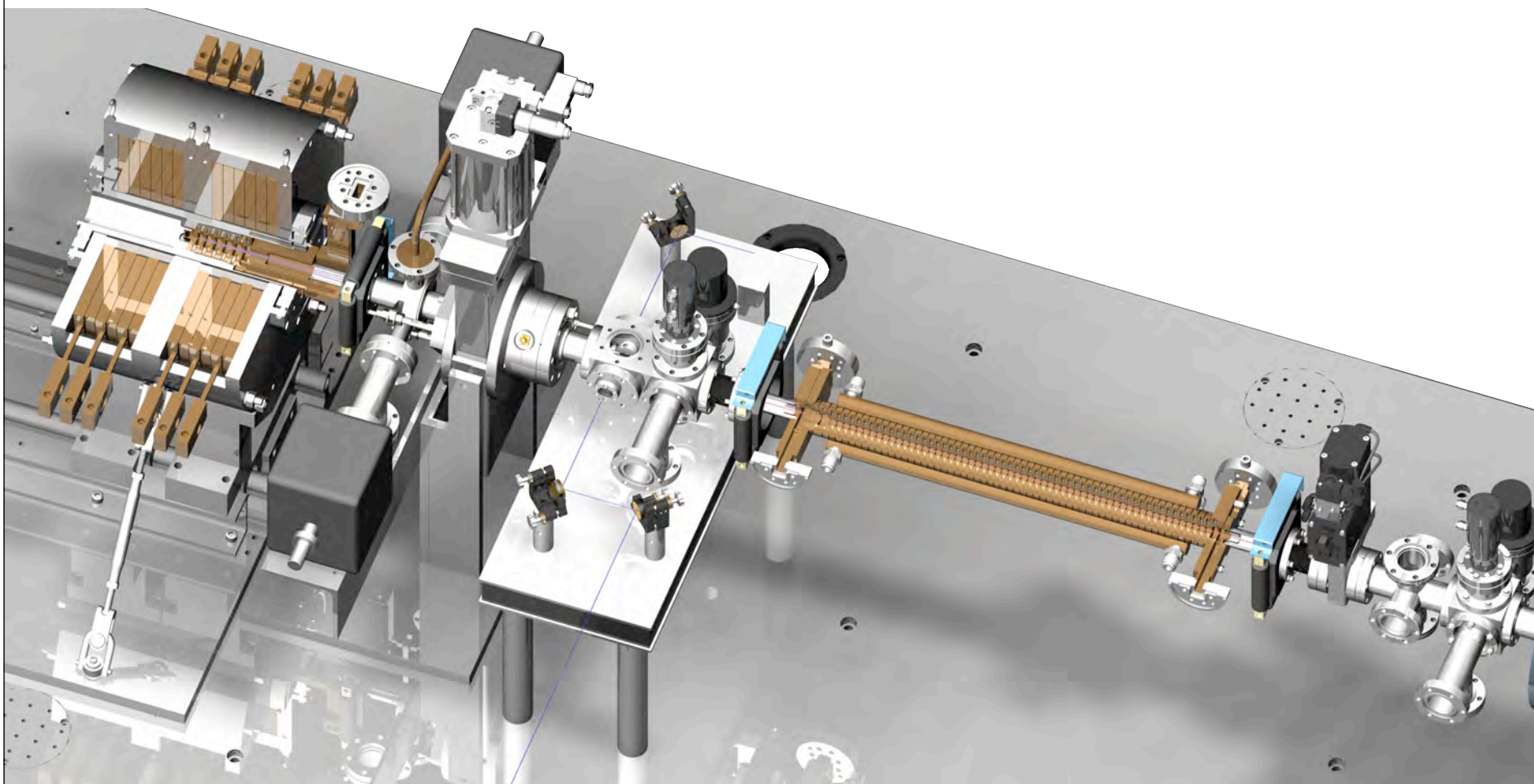


Worldwide high power x-band sources



LLNL set up is currently the only facility where high quality x-band RF is coupled with state-of-the-art structures to produce beam

Photo-gun and first section



X-band photo-gun evolution



Mark 0

200 MV/m

**SLAC 5.5 cell x-band gun
designed by Arnold Vlieks (SLAC)**

Mark 1

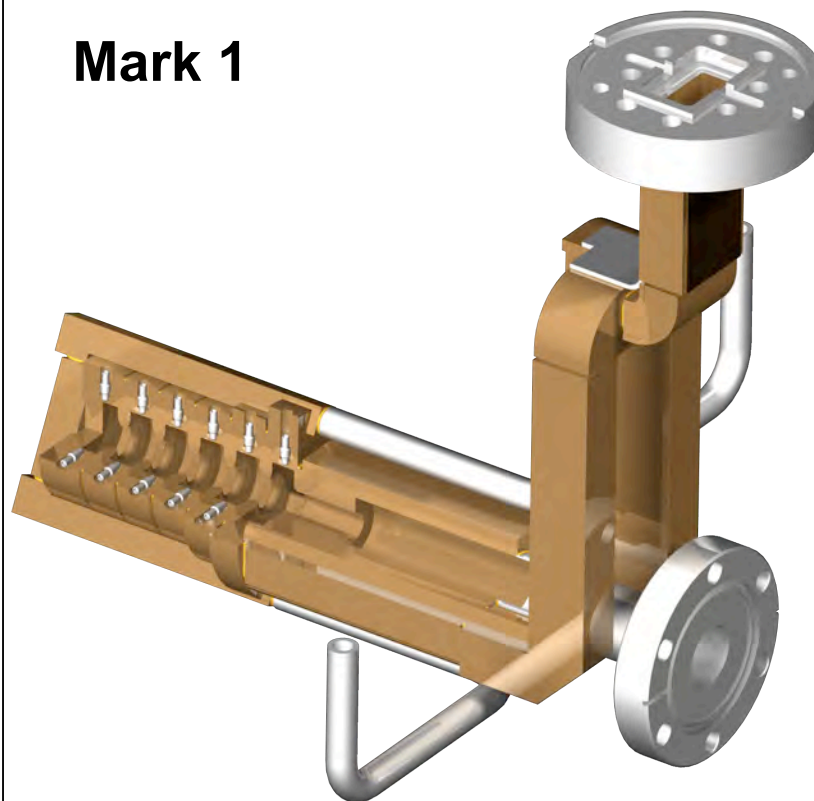
**LLNL/SLAC 5.59 cell x-band gun
design lead - Roark Marsh (LLNL)**

X-band photo-gun evolution



- Longer Half cell for lower final emittance
- Better mode separation for less mode beating on cathode surface
- Elliptical irises for lower peak surface electric field
- Dual feed racetrack coupler for minimized RF quadrupole kick
- Optimized beta for a balance of fast gun fill time and low pulsed heating

Mark 1



**LLNL/SLAC 5.59 cell x-band gun
design lead - Roark Marsh (LLNL)**

Multiple Codes were used to design and verify the Mark 1 x-band RF photo-gun

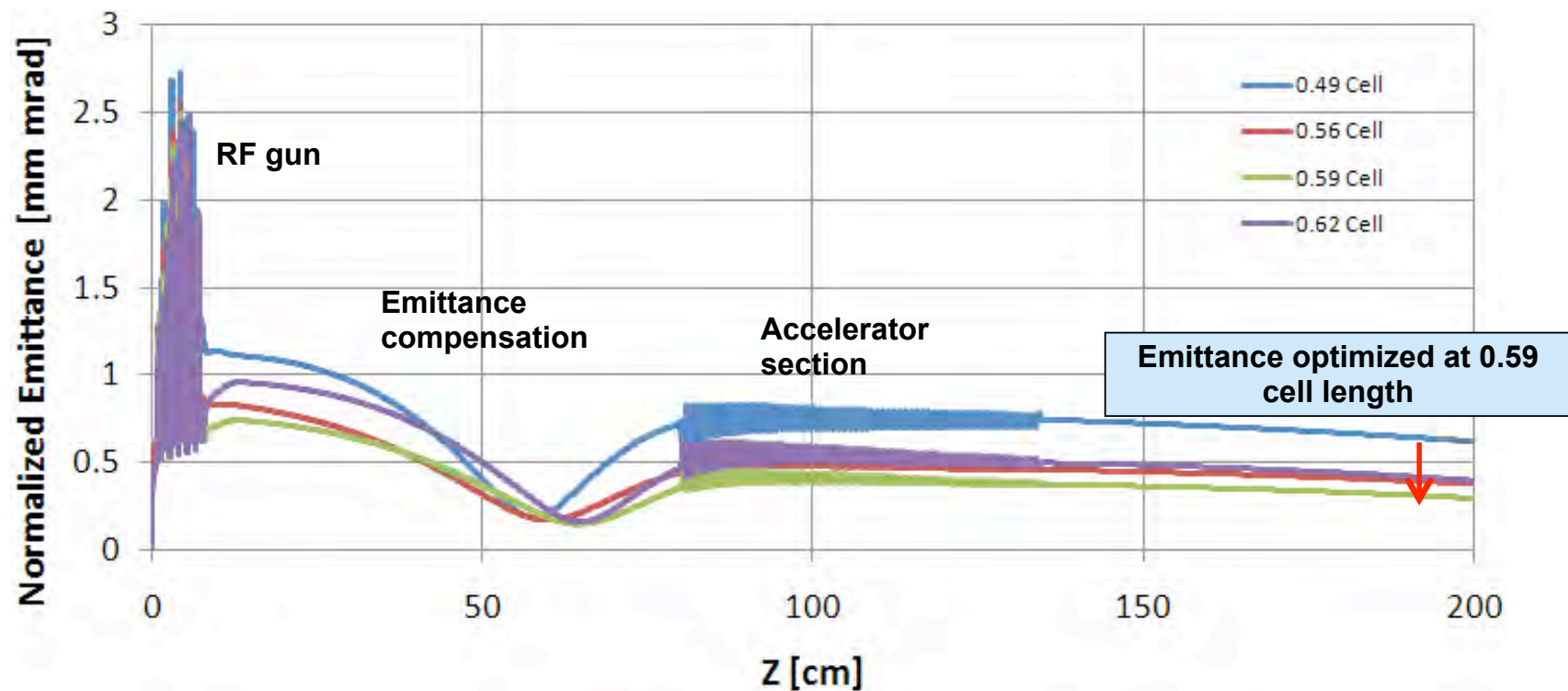


- **Pro/Engineer and Solidworks**
 - Mechanical design
- **Superfish**
 - Fast 2D axially symmetric gun optimization
- **PARMELA, GPT, ASTRA, Impact-T**
 - Beam dynamics
- **Ansoft HFSS**
 - Full 3D microwave design, frequency domain
- **ACE3P**
 - Suite of 3D parallel SLAC codes
 - CUBIT mesher, Omega3P, S3P, T3P, PIC3P, ParaView postprocessing
 - Final benchmark and verification

Redesigned longer half cell for optimized brightness



Optimized launch phase and solenoid strength
Beam parameters: $Q = 250$ pC, $\tau\phi = 10$ deg.
200 MV/m cathode field

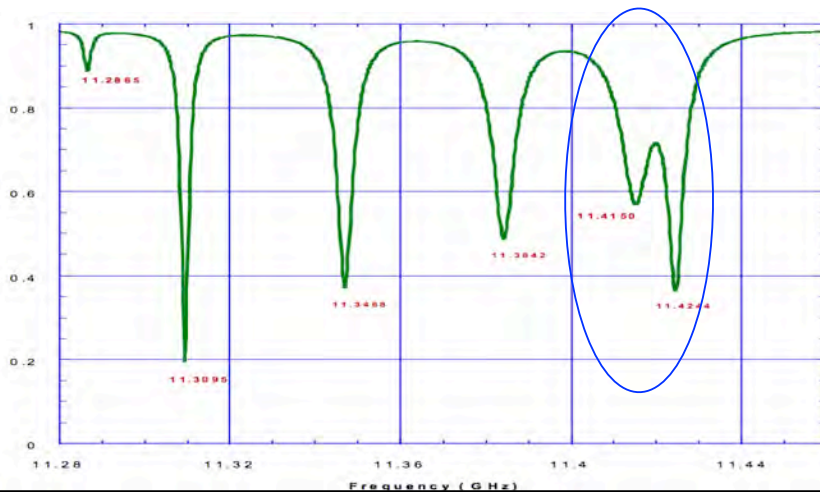


Calculated mode separation is greater than Mark 0



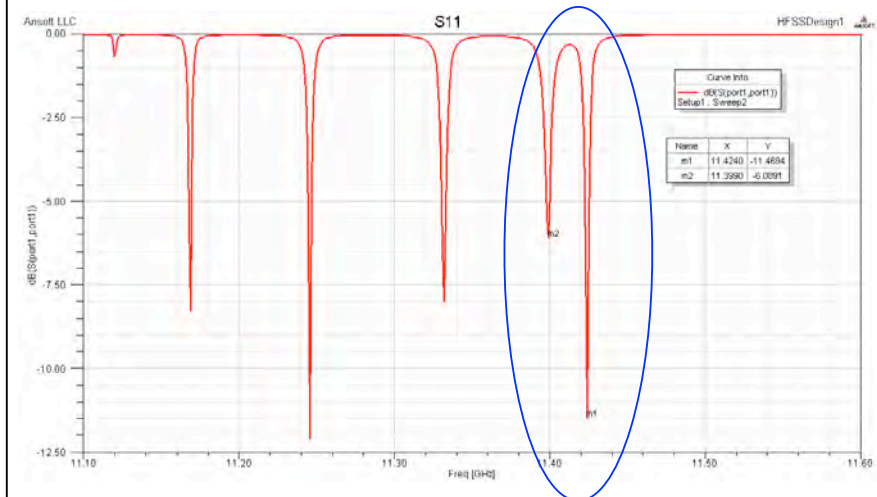
- Original Mark 0 design provided <10 MHz separation
- Redesign of iris geometry achieves 25 MHz spacing

Mark 0 Measurement



9 MHz separation: Clear co-excitation

Mark 1 Design

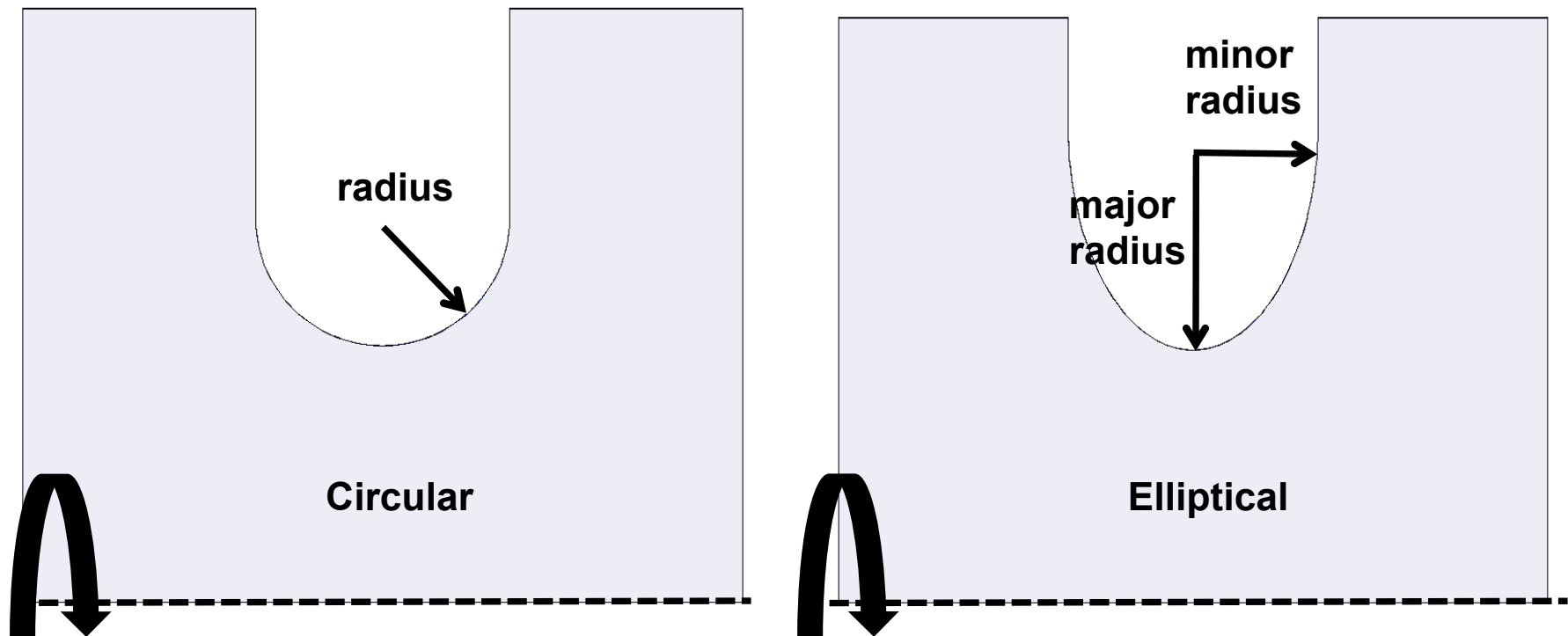


25 MHz separation: Distinct modes

Iris geometry was changed from circular to elliptical and optimized



- Ellipticity for Mark 1 RF gun: major/minor = 1.8

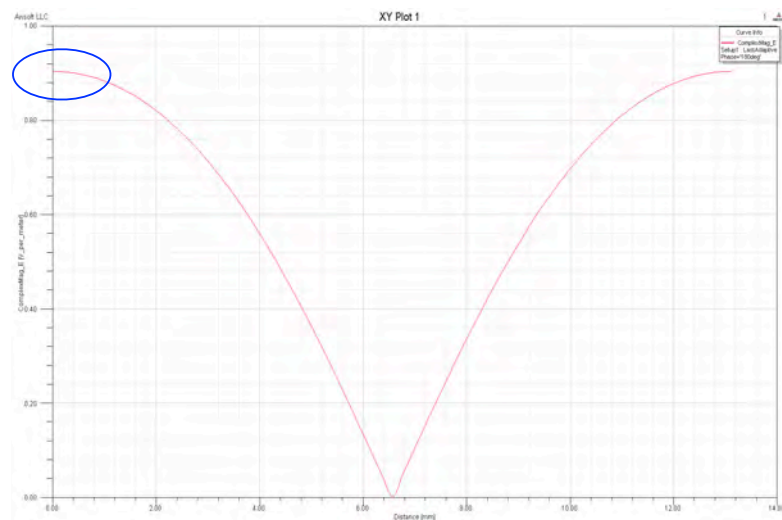


Electrical performance difference



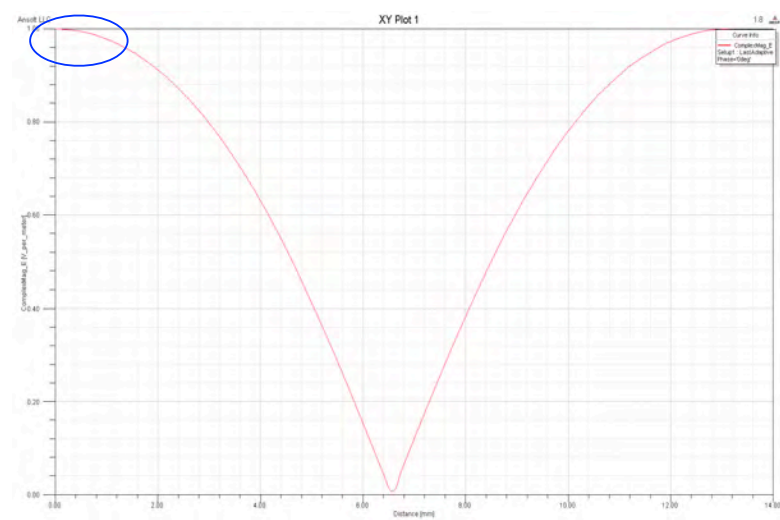
- Figure of merit: ratio of peak axial electric field to peak surface electric field on iris

Axial Field for Circular Iris



Ratio of 0.92

Axial Field for Elliptical Iris



Ratio of 1.04

Peak axial field improvement of >10% for the same peak surface electric field

B194 X-Band Test Station

PHYSICAL REVIEW SPECIAL TOPICS - ACCELERATORS AND BEAMS, VOLUME 00, 0

Modeling and design of an X-band rf photoinjector

R. A. Marsh,* F. Albert, S. G. Anderson, G. Beer, T. S. Chu, R. R. Cross, G. A. Deis, C. A. Ebbers, D. J. Gibson, T. L. Houck, F. V. Hartemann, and C. P. J. Barty

Lawrence Livermore National Laboratory, Livermore, California 94550, USA

A. Candel, E. N. Jongewaard, Z. Li, C. Limborg-Deprey, A. E. Vlieks, F. Wang, J. W. Wang, F. Zhou, C. Adolphsen, and T. O. Raubenheimer

SLAC National Accelerator Laboratory, Menlo Park, California 94025, USA

(Received 28 March 2012)

A design for an X-band rf photoinjector that was developed jointly by Stanford Linear Accelerator Center (SLAC) and Lawrence Livermore National Laboratory (LLNL) is presented. The photoinjector is based around a 5.59 cell rf gun that has state-of-the-art features including: elliptical contoured trises; improved mode separation; an optimized initial half cell length; a racetrack input coupler; and coupling that balances pulsed heating with cavity fill time. Radio-frequency and beam dynamics modeling have been done using a combination of codes including PARMELA, HFSS, IMPACT-T, ASTRA, and the ACEP suite of codes developed at SLAC. The impact of lower gradient operation, magnet misalignment, solenoid multipole errors, beam offset, mode beating, wakefields, and beam line symmetry have been analyzed and are described. Fabrication and testing plans at both LLNL and SLAC are discussed.

DOI:

PACS numbers: 07.77.Ka, 29.25.Bx, 29.27.-a, 41.75.Hr

I. INTRODUCTION

The development of rf photoinjector technology has enabled free electron lasers and other fourth generation light sources, such as the Linac Coherent Light Source (LCLS) at Stanford Linear Accelerator Center (SLAC) National Accelerator Laboratory. At Lawrence Livermore National Laboratory (LLNL) a novel gamma-ray light source is being developed and built taking advantage of inverse Compton scattering to extract energy from high brightness electron bunches to boost laser photons to MeV energies. At SLAC advanced compact X-band (11.424 GHz) photoinjector R&D is being done to investigate the possibility of generating short enough high brightness electron bunches that at least one stage of bunch compression may become unnecessary, as well as to enable an all X-band free-electron laser by establishing a proven electron beam source.

Nuclear resonance fluorescence (NRF) is a process in which a nucleus, excited by gamma rays, reradiates high-

the only process capable of producing a narrow bandwidth radiation (below 1% $\Delta\omega/\omega$) at gamma-ray energies by using state-of-the-art accelerator and laser technologies. In Compton scattering sources, a short laser pulse and a relativistic electron beam collide to yield tunable, monochromatic, polarized gamma-ray photons. Building on prior work on narrowband gamma-ray light sources at LLNL [1–8], the LLNL Nuclear Photonics Facility (NPF) will be equipped with a tunable MEGA-ray source using an all X-band linac including an X-band rf photoinjector. The advantages of operating at X-band and further detail on the linac design are available in [9].

This paper describes an rf photoinjector which will be tested at the X-band test area (XTA) at SLAC, at an X-band test station at LLNL, and will serve as the injector for the X-band very energetic light for the observation and characterization of isotopic resonances and the assay and precision tomography of objects with radiation (VELOCIRAPTOR) linac, designed to drive the precision,

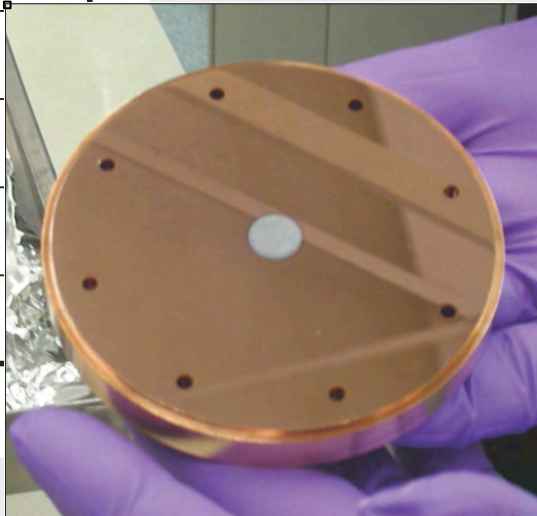
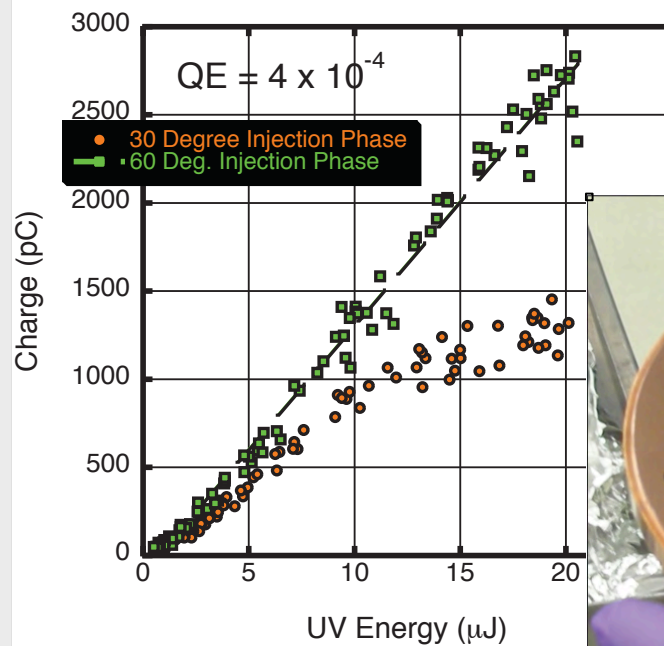
WARNING
Low Voltage/High
Current: Treat
Exposed Metal as
Energized

High efficiency cathode materials reduce the requirements for the photo-gun drive laser



Mg QE Demonstration

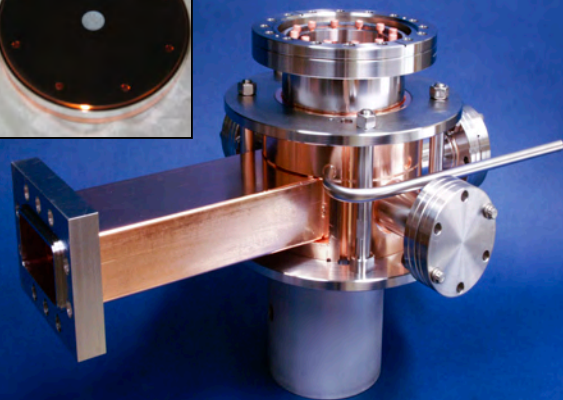
Charge Extracted from Mg Photo-cathode
Gun Peak Field = 120 MV/m



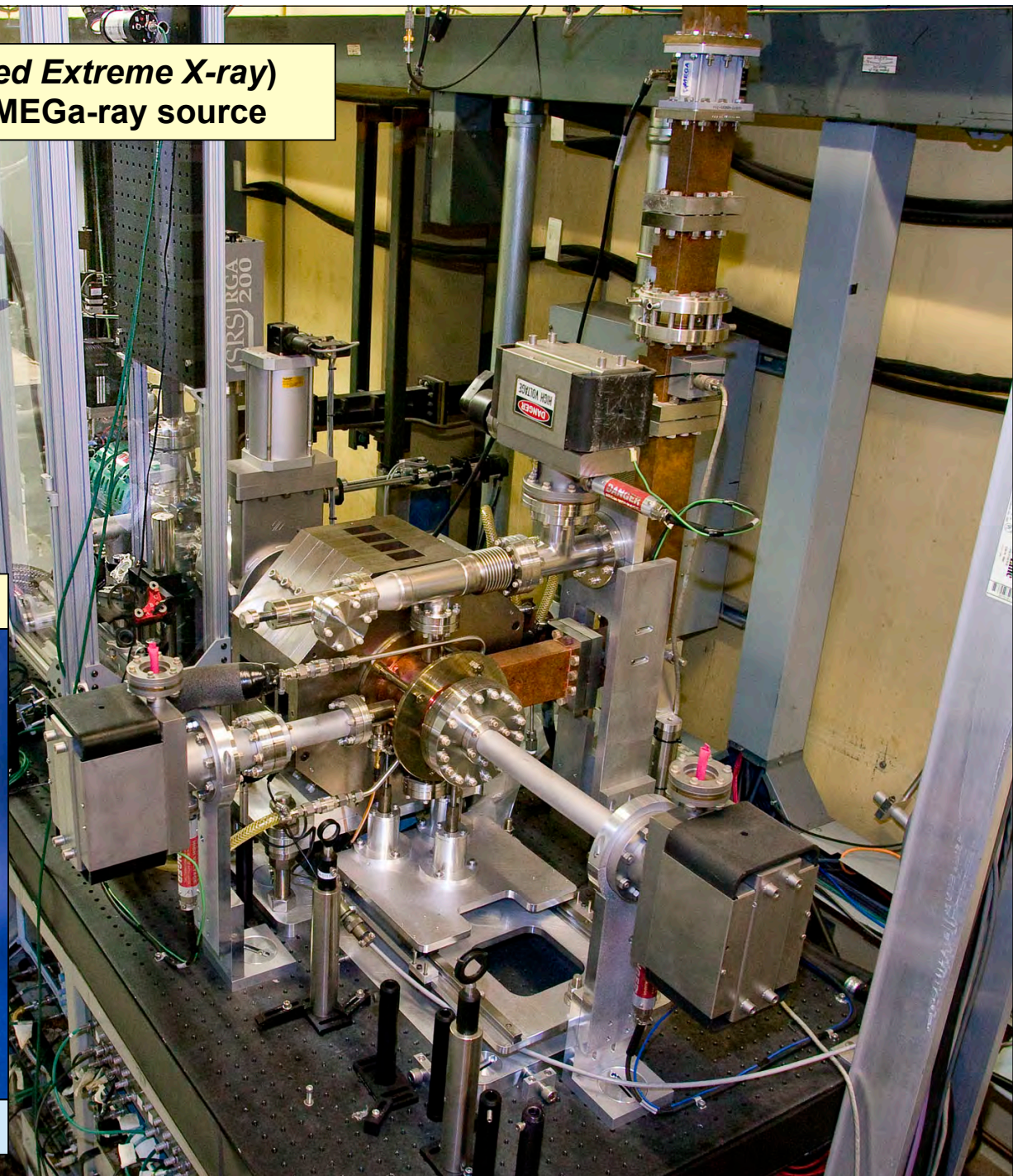
10^{-4} QE demonstrated in S-Band gun

The T-REX (*Thomson-Radiated Extreme X-ray*) project created LLNL's first MEGa-ray source

Ultra-low Emittance Photogun

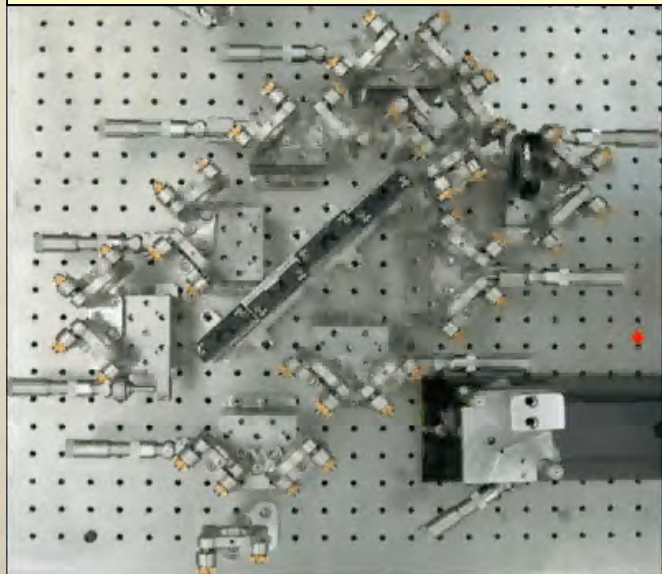


Fully Symmetrized w/ Mg Cathode



**The T-REX (*Thomson-Radiated Extreme X-ray*)
project created LLNL's first MEGa-ray source**

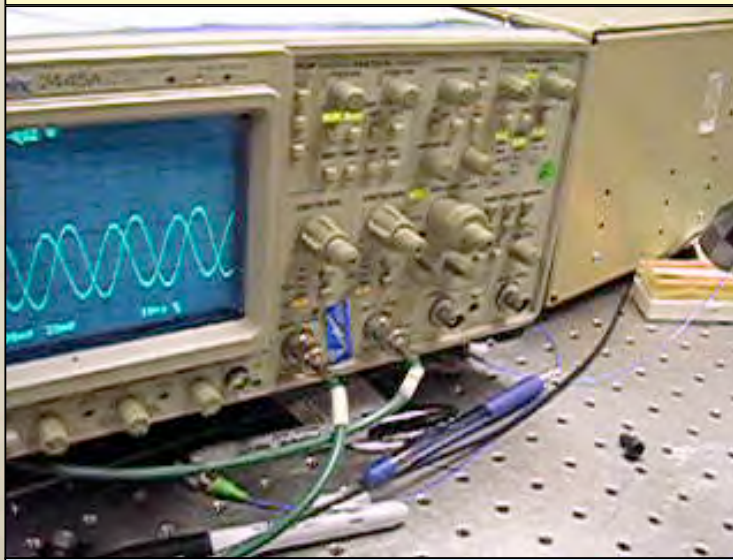
Hyper-Michelson Pulse Shaper



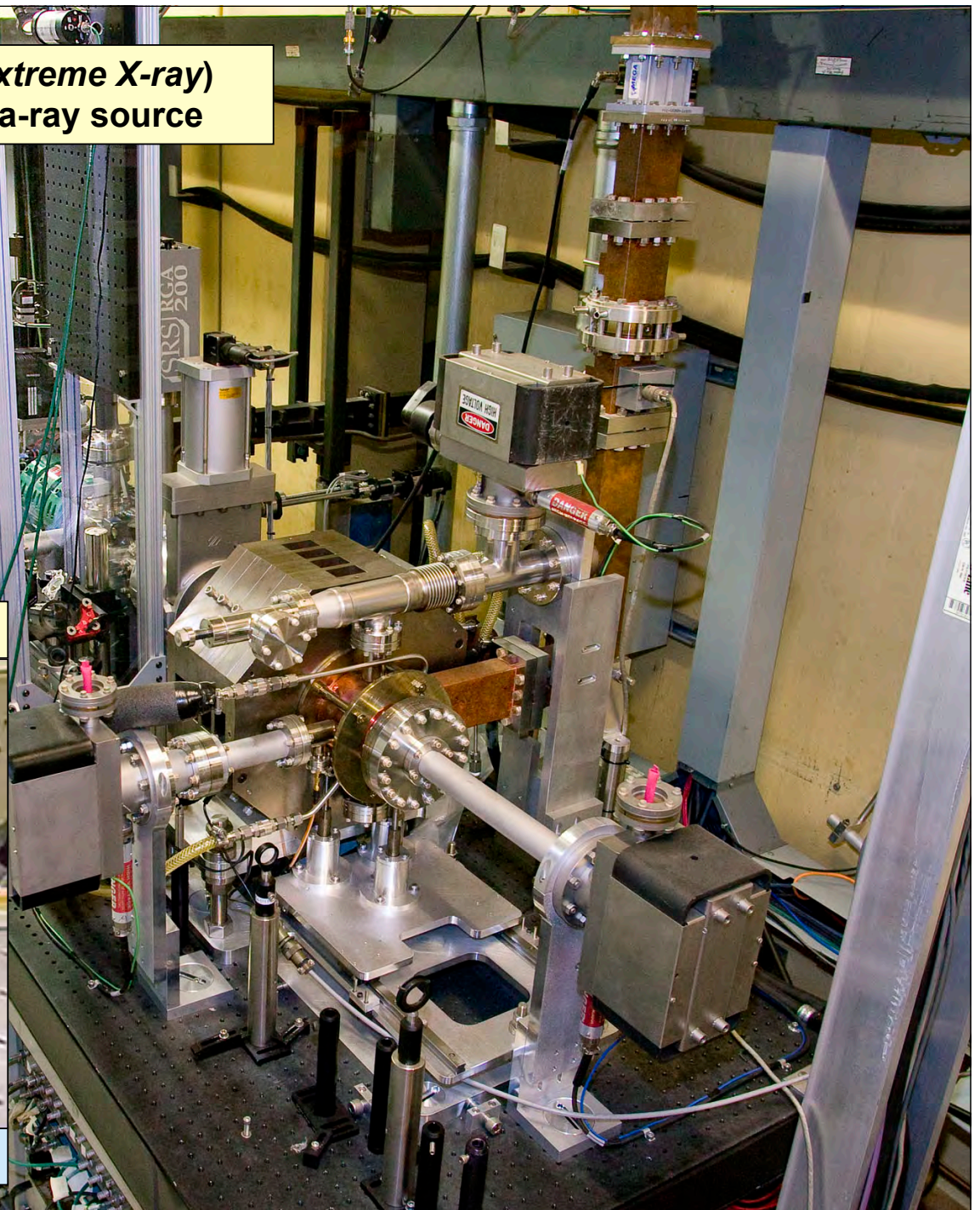
100 fs rise time, 10 ps flat top pulses

**The T-REX (*Thomson-Radiated Extreme X-ray*)
project created LLNL's first MEGa-ray source**

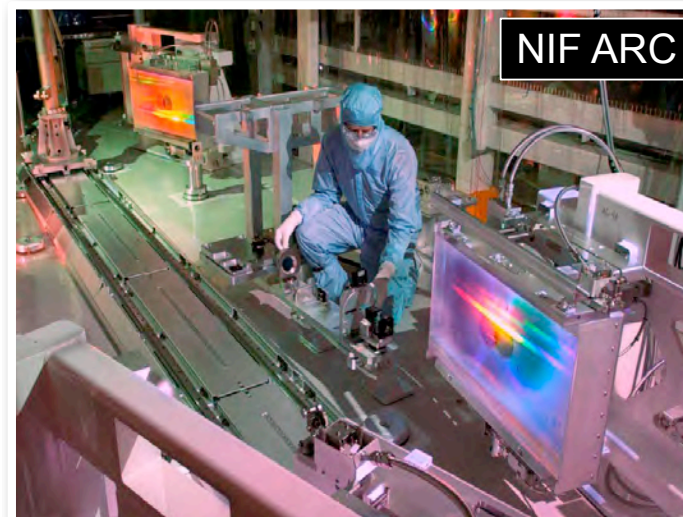
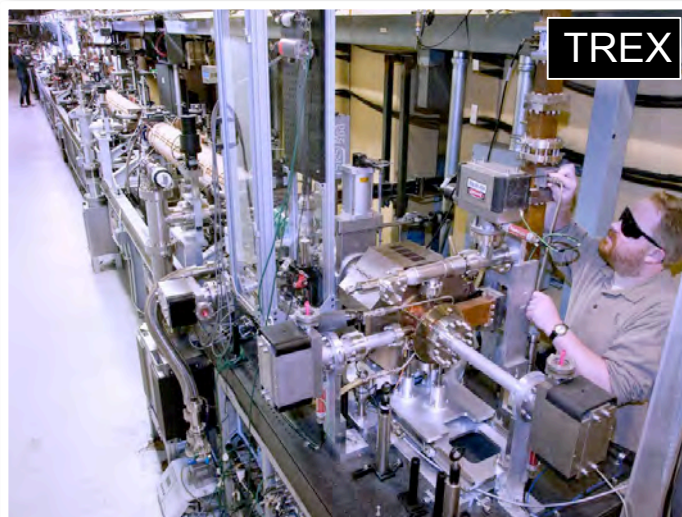
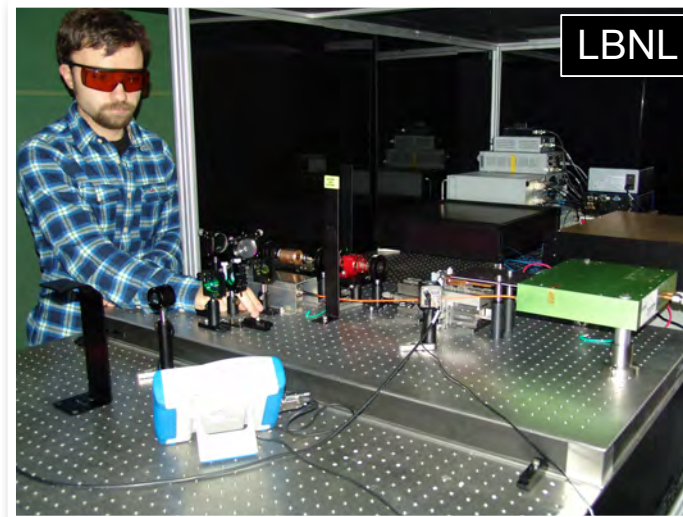
Robust Fiber Front end



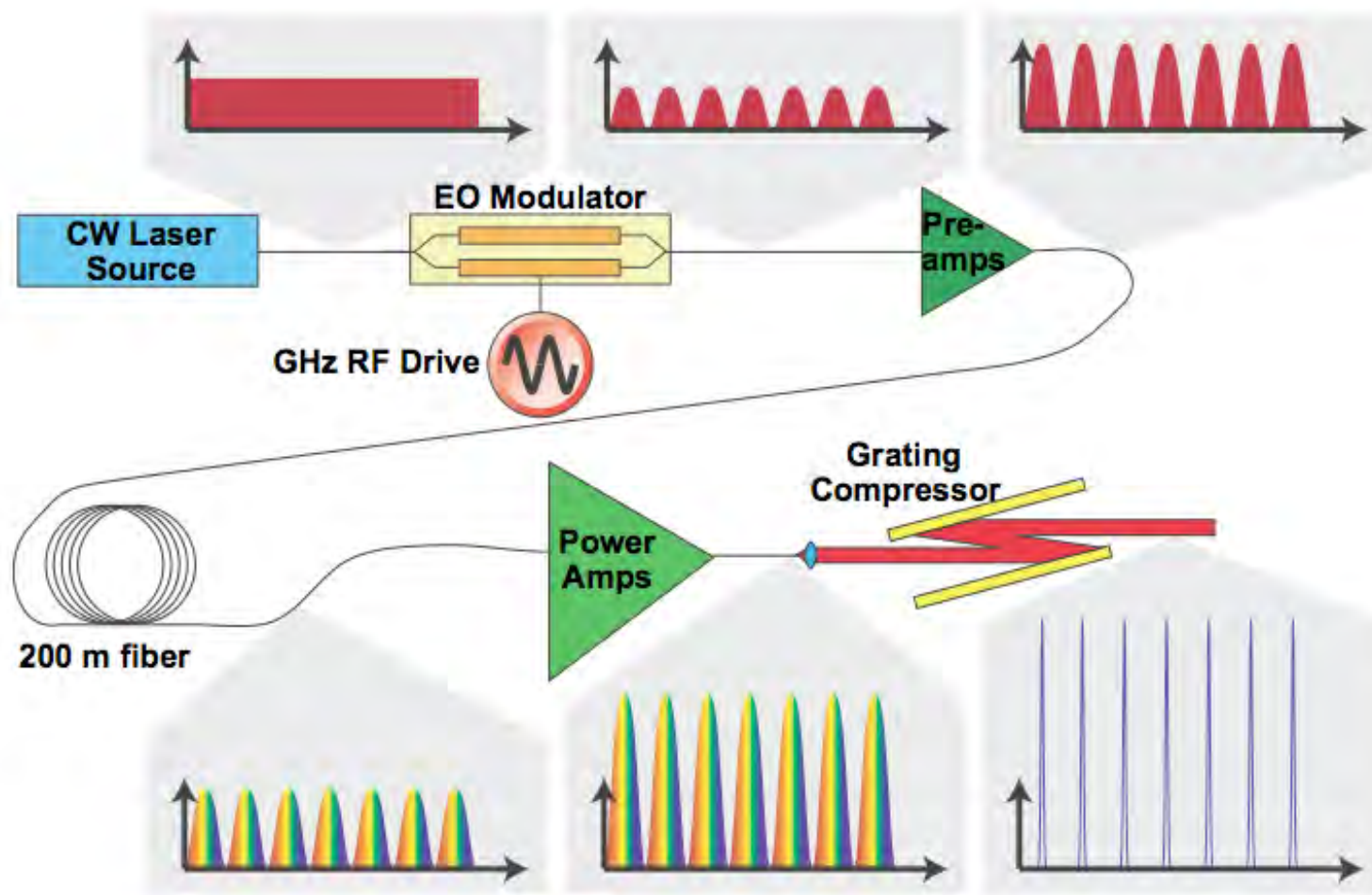
<100 fs pulses locked to S-band RF



Fiber lasers at LLNL



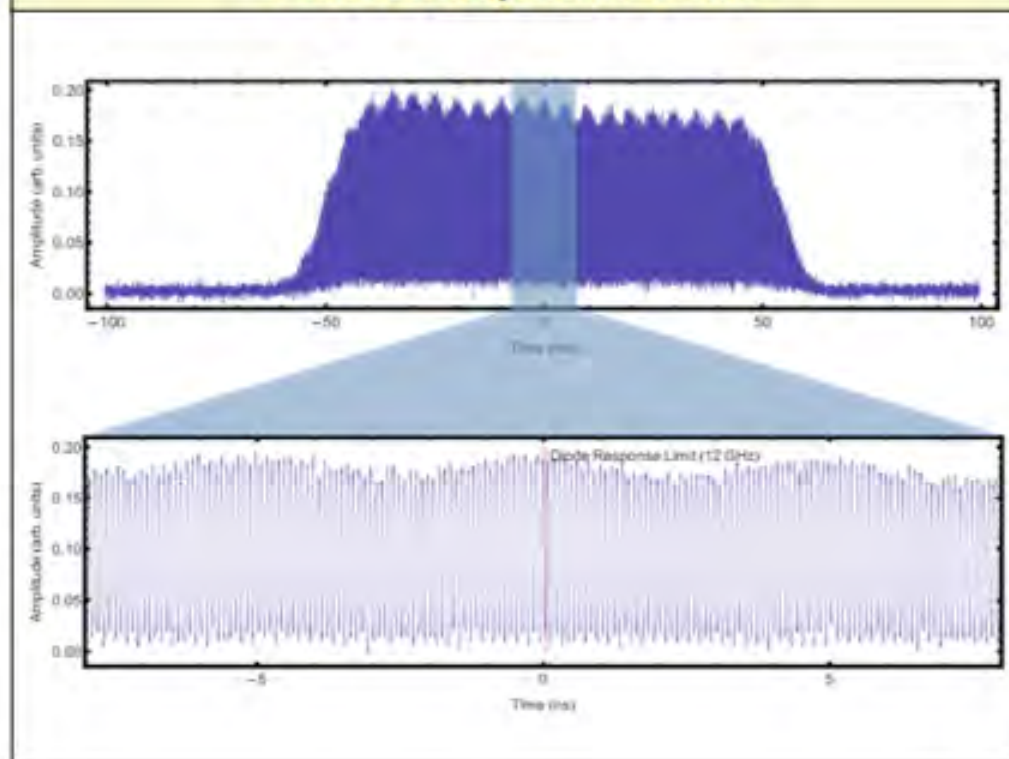
“CW” method for generation of 11.424 GHz, synchronized train of picosecond IR pulses



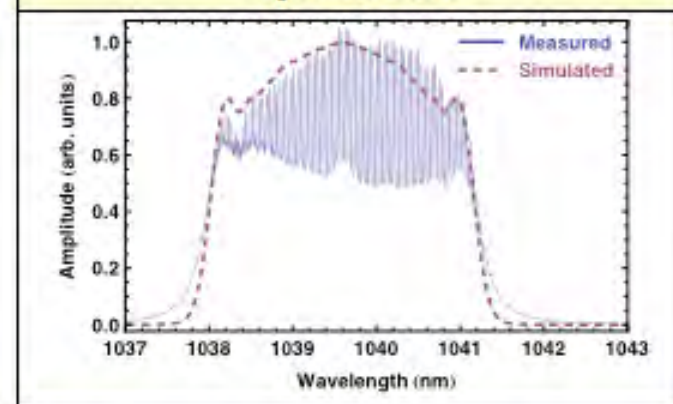
“CW” method for generation of 11.424 GHz, synchronized train of picosecond IR pulses



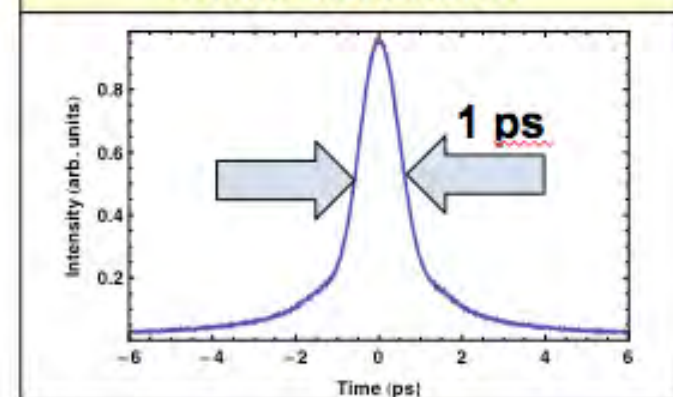
Burst Temporal Profile



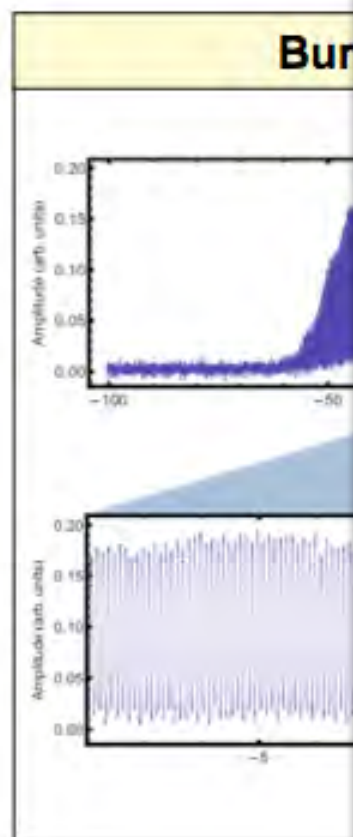
Spectrum



Autocorrelation



“CW” method for generation of 11.424 GHz, synchronized train of picosecond IR pulses



Widely tunable 11 GHz femtosecond fiber laser based on a non-modelocked source

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An 11 GHz fiber laser built on a modulated CW platform is described and characterized. This compact, vibration-insensitive, fiber based system can be operated at wavelengths compatible with high energy fiber technology, is driven by an RF signal directly, and is tunable over a wide range of drive frequencies. The demonstration system when operated at 1040 nm is capable of 50 ns bursts of 575 micro-pulses produced at a macro-pulse rate of 83 kHz where the macro-pulse and micro-pulse energies are 1.8 μ J and 3.2 nJ respectively. Its micro-pulses are compressed to a duration of 850 fs.

OCIS Codes: 060.3510, 140.7090

Very-short-pulse laser sources (in the range of picoseconds to femtoseconds) with high repetition rates (100 MHz to tens of GHz) are needed to drive short-wavelength high-energy photon sources via higher-order nonlinear optical parametric interactions, and as photocathode illuminators to create photo-electrons in high frequency particle accelerators. Other applications of these ultrafast pulse laser systems include materials processing, 3-D lithography, high-data-rate laser communication, and remote sensing systems. The system architecture described here is wavelength compatible with high energy fiber technology, is driven directly by an RF source (and thus sidesteps synchronization issues), and allows for amplification to materials-processing pulse energies. Moreover, the ability to modify electronically the temporal pulse shape, in amplitude and phase [1, 2], offers the possibility of controlling various complex photochemical processes and quantum control of interactions on molecular time scales.

We generate a laser pulse train by modulating a continuous-wave (cw) laser with an RF source. Several groups [2, 3] have converted cw lasers to sub-ps, high frequency pulse trains; however, these groups have relied on “time-lens” techniques [4] to generate ps-level bandwidths, then used soliton compression at 1550 nm in specially optimized fibers to generate further bandwidth while simultaneously compressing the pulse. At the 1030-1070 nm wavelengths where we wish to work, the soliton compression scheme is not feasible because the dispersion in standard fibers has the opposite sign from dispersion at 1550 nm. In the demonstration of this cw-modulation concept reported here, we rely on self-phase modulation (SPM) [5] to generate 3.2 nm of bandwidth and compress the pulse with a grating-pair compressor. To reduce allow high pulse energies at modest average laser power, we

the groups (macro-pulses) occur at multi-kHz rates while the individual pulses (micro-pulses) are at multi-GHz rates. Alternatively, the micro-pulses might be directly amplified to modest energies for a clock-distribution scheme, or driven to micro-joule energies to create tens of kilowatts of power for machining or supercontinuum sources.

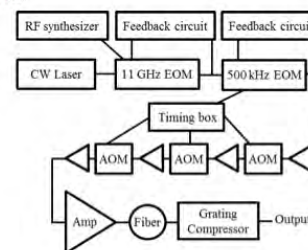
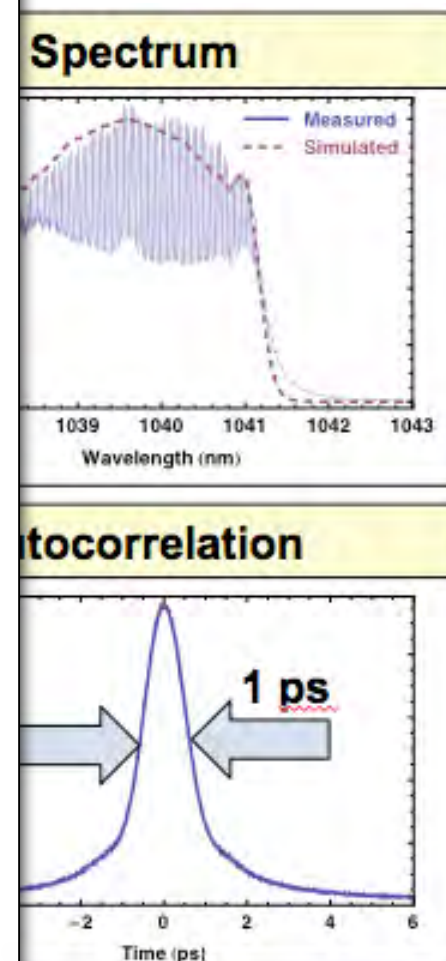
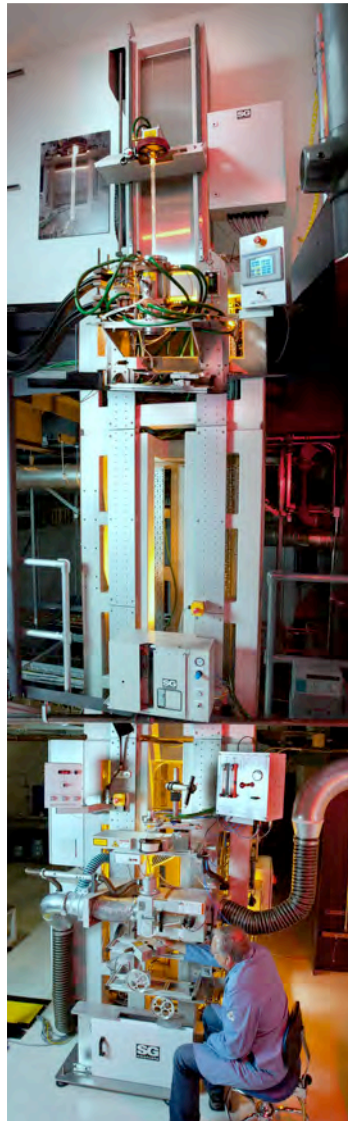


Figure 1 Schematic diagram of the experiment.

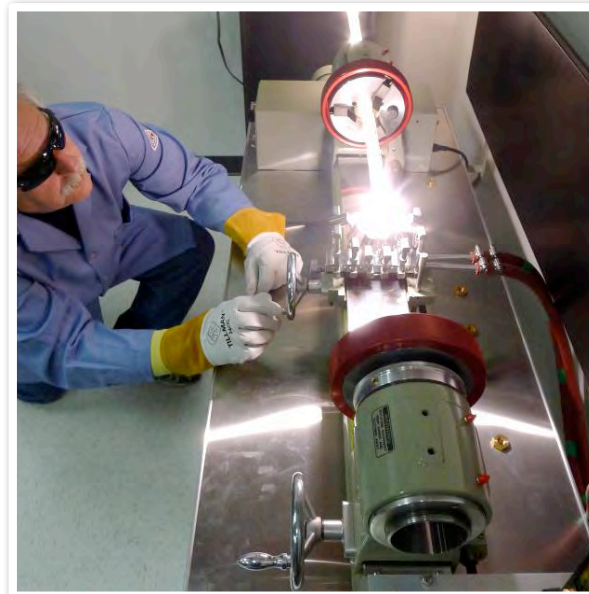
The architecture, shown in Figure 1, begins with a cw laser: a New Focus Velocity laser set to provide a 1040 nm beam. The beam is sent through an EOSPACE-brand, Z-cut, 20 GHz, dual drive Mach-Zehnder electro-optic modulator (EOM) monitored by a control circuit (YY Labs, Inc.) to keep the modulator null-biased – that is, biased to block light when the RF drive is off. This modulator is driven with 20.1 dBm of 5.7 GHz RF power. Because of the null bias, the RF creates an 11.4 GHz laser pulse train with 44 ps pulse length and no cw component as shown in Figure 2: the latter prevents stimulated Brillouin scattering from damaging subsequent fiber amplifiers. A second EOM temporally slices macro-pulses at a 500 kHz rate with 50 ns duration; thus, each macro-pulse contains approximately



LLNL fiber facilities



- Fiber fabrication
 - 8.2m draw tower
 - Preform assembly fixtures
 - Glass working lathe

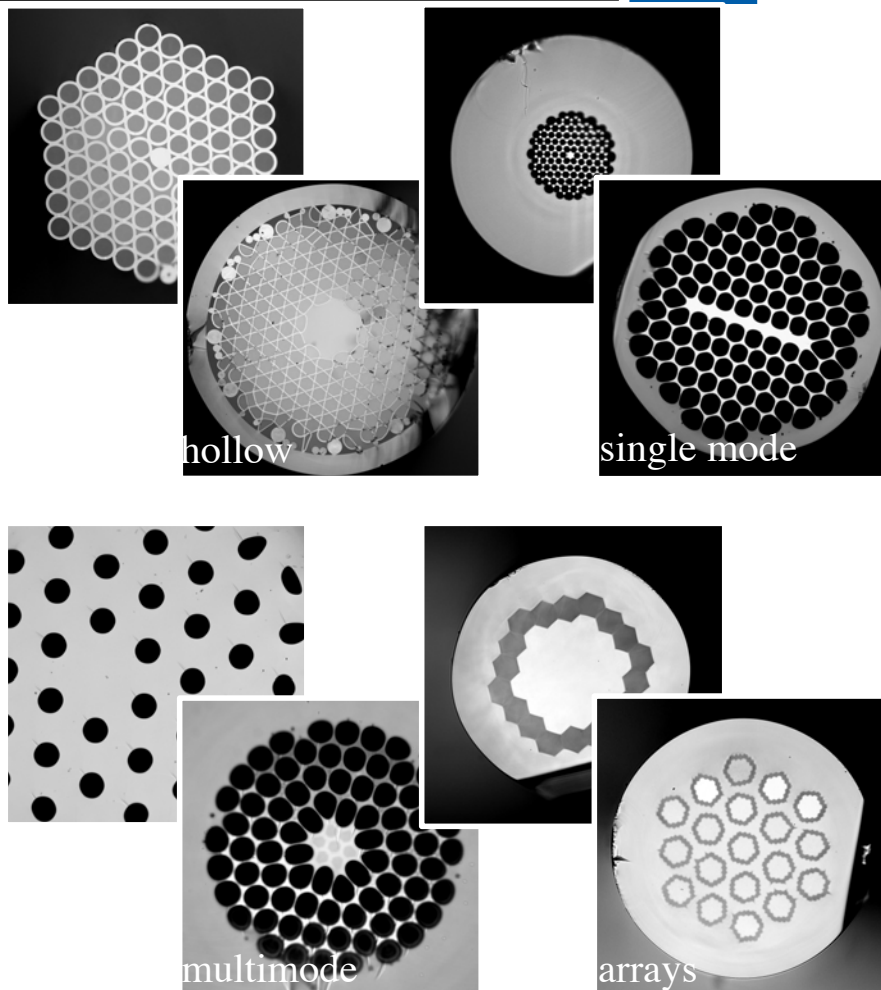


Recent work



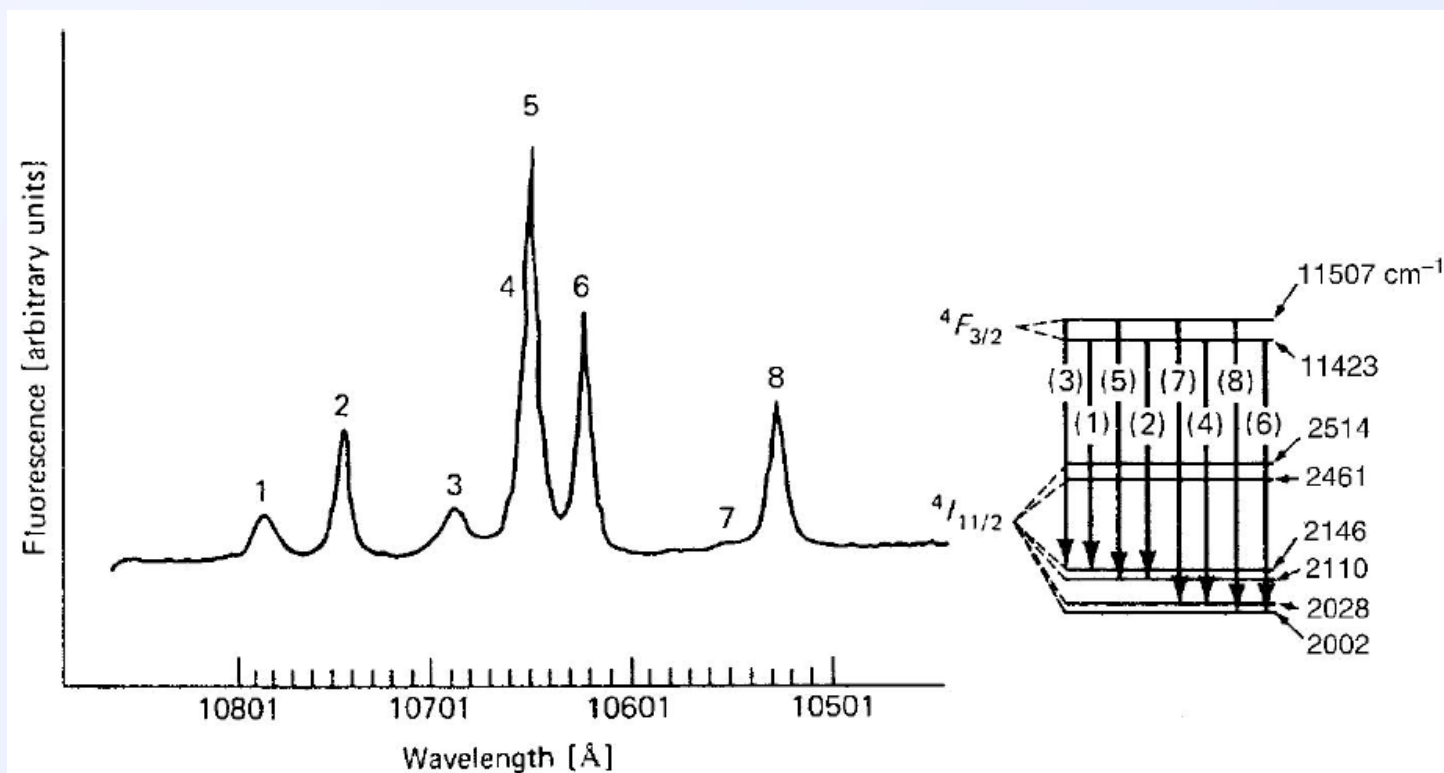
- Discerned laser limits¹
- Fabricated ribbon fiber²
- Launched high-order modes³
- Built, tested lasers⁴
- Designed new type of fiber^{5,6,7}

1. J. Dawson et al, Opt. Exp. 16 13240-13266 (2008).
2. D. Drachenberg et al, Opt. Exp. 21 11257-11269 (2013).
3. A. Sridharan et al, Opt. Exp. 20 28792-28800 (2012).
4. D. Drachenberg et al, submitted to Opt. Express (2013).
5. M. Messerly et al, Opt. Exp. 21 12683-12698 (2013).
6. M. Messerly et al, Optics Letters 38 3329-3332 (2013).
7. P. Pax et al, submitted to Optics Letters (2013).



Nd:YAG could be ideal for a drive laser

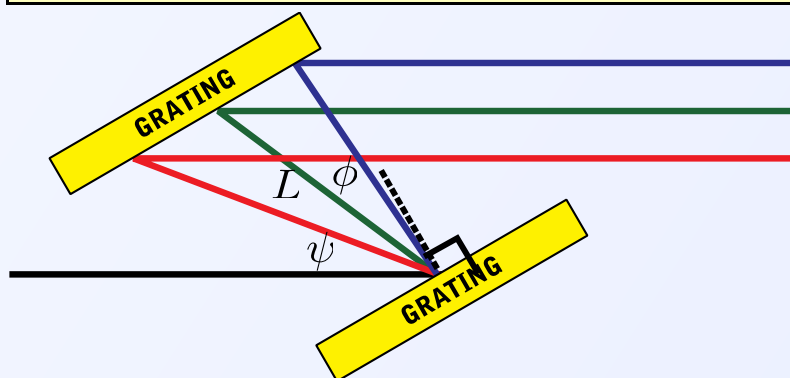
- Routinely used to produce few-ns long multi-Joule laser pulses
- Sufficient bandwidth for ps pulses: 120 GHz (~ 0.5 nm)
- **Requires stretching < 1 nm to a few ns**



Koechner and Bass, *Solid-State Lasers*, Springer-Verlag, New York (2003).

CPA with narrow-band pulses requires very large dispersion

Standard Compressor



$$\frac{d\tau}{d\lambda} = \frac{2L\lambda}{cd^2 \cos^2 \phi} = 6000 \frac{\text{ps}}{\text{nm}}$$

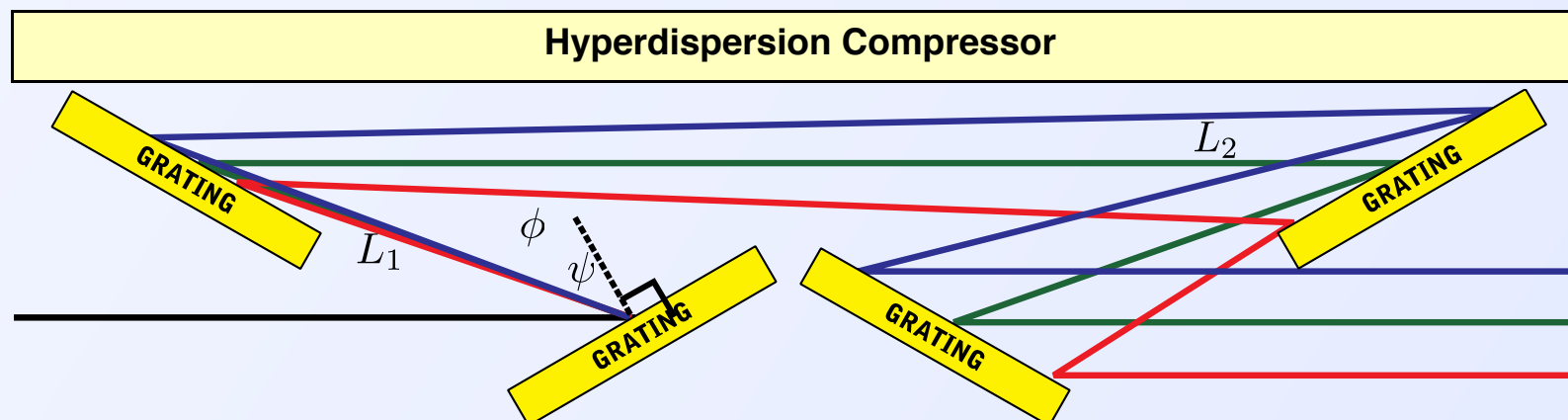
$$\frac{1}{d} = 1740 \frac{\text{lines}}{\text{mm}}$$

$$\psi = 65^\circ$$

$$\lambda = 1064 \text{ nm}$$

$$L \approx 30 \text{ m}$$

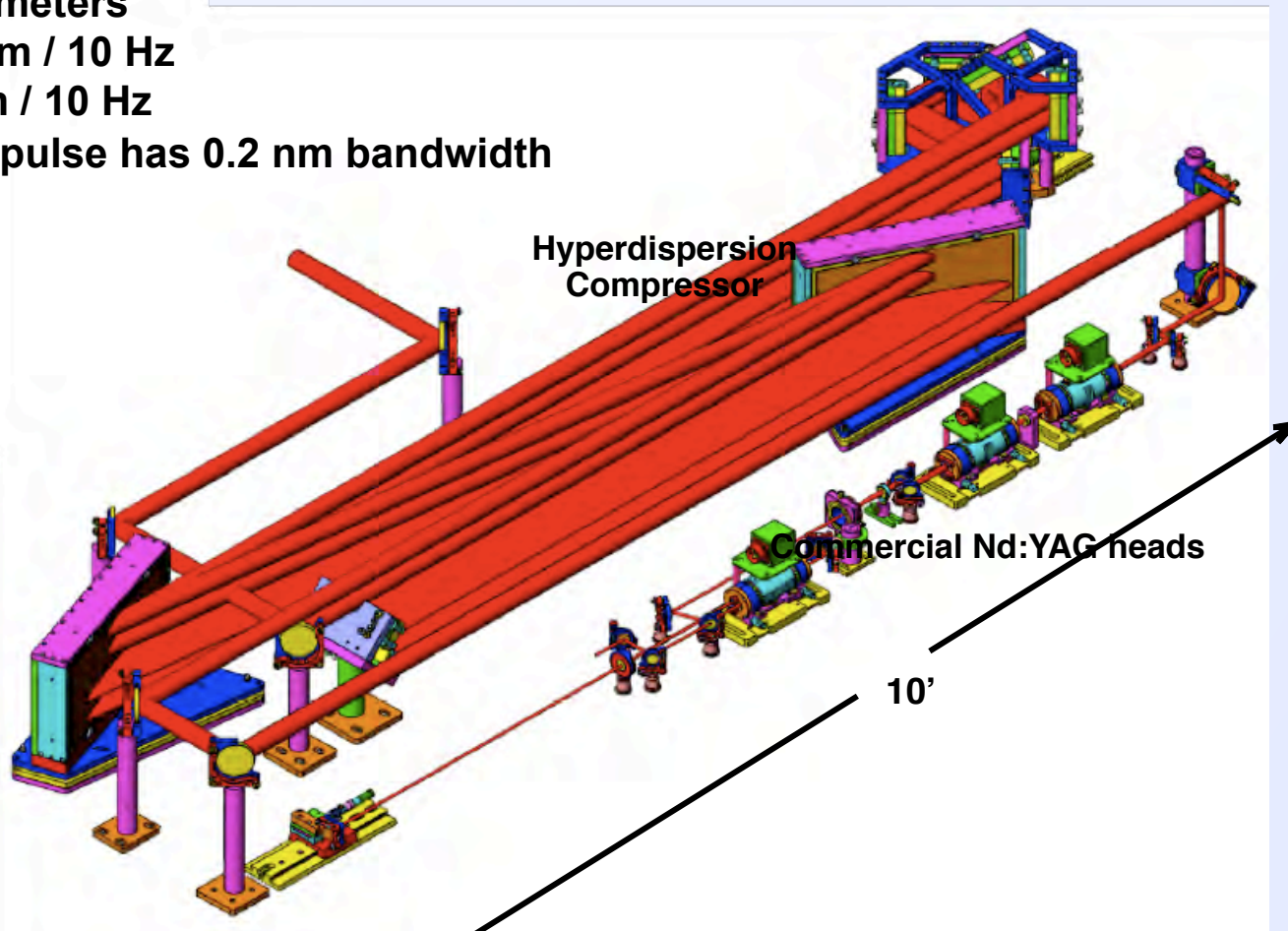
Hyperdispersion Compressor



T-REX Interaction Laser demonstrated Hyperdispersion Stretcher/Compressor

- Commercial flashlamp-pumped heads
- Seeded with fiber system similar to PDL front end
- Laser System Parameters
 - 0.8 J / 1064 nm / 10 Hz
 - 0.3 J / 532 nm / 10 Hz
- Amplified 1064 nm pulse has 0.2 nm bandwidth

Three Continuum
Amplifier Heads



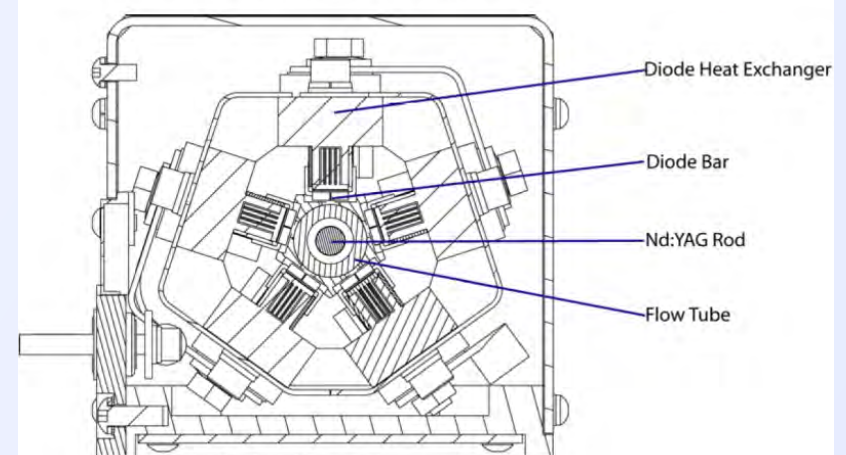
Commercial diode pumped laser heads are now capable of several 100 W cw operation



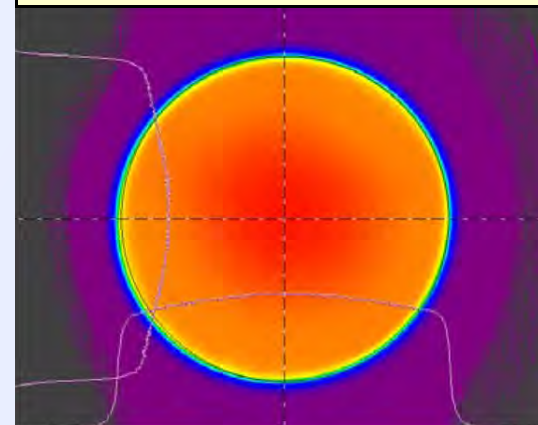
Diode pumped amplifier head w/ 1 cm diameter x 14.6 cm long Nd:YAG rod



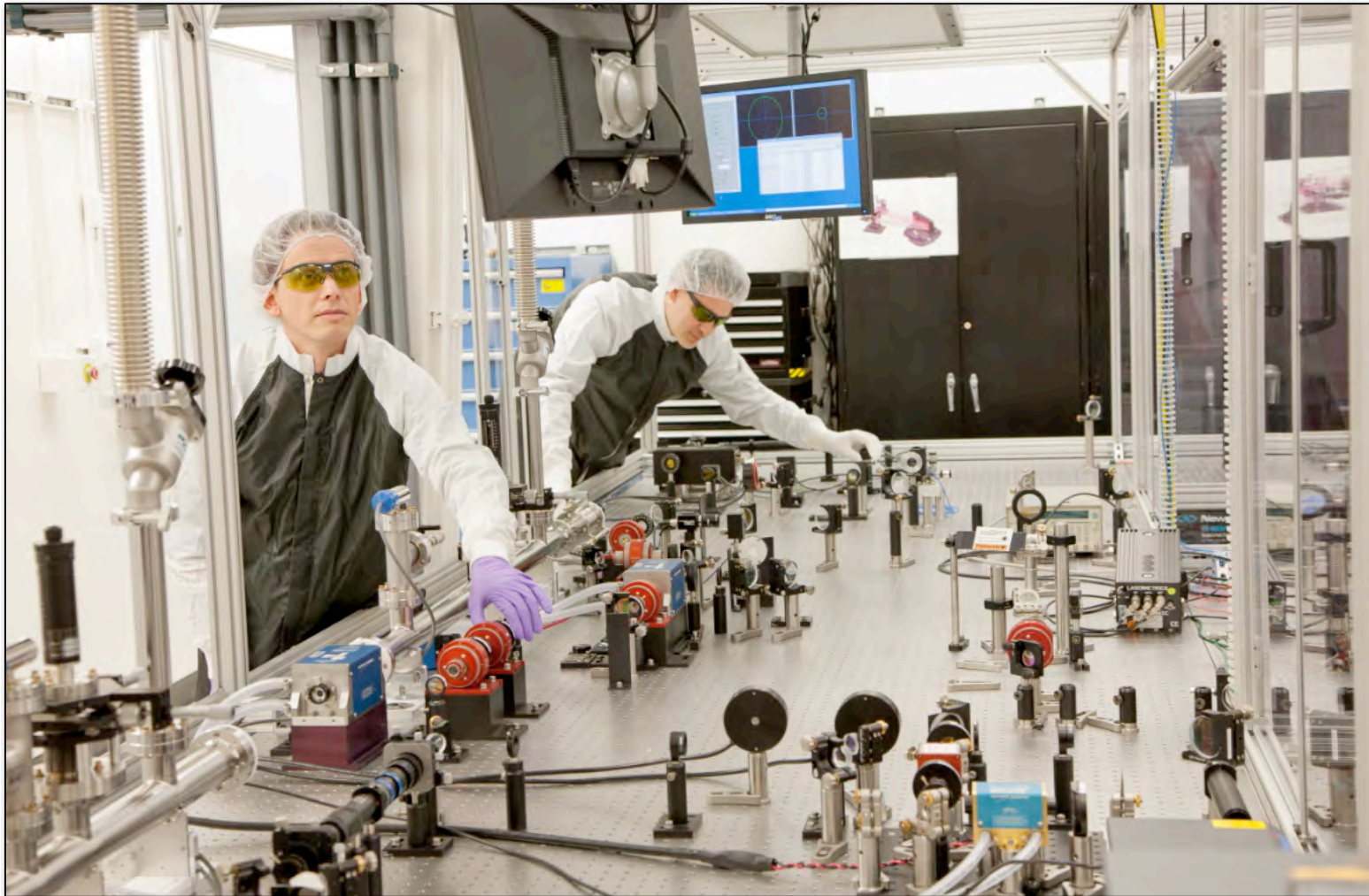
In pulsed mode these heads should be capable of joule level operation



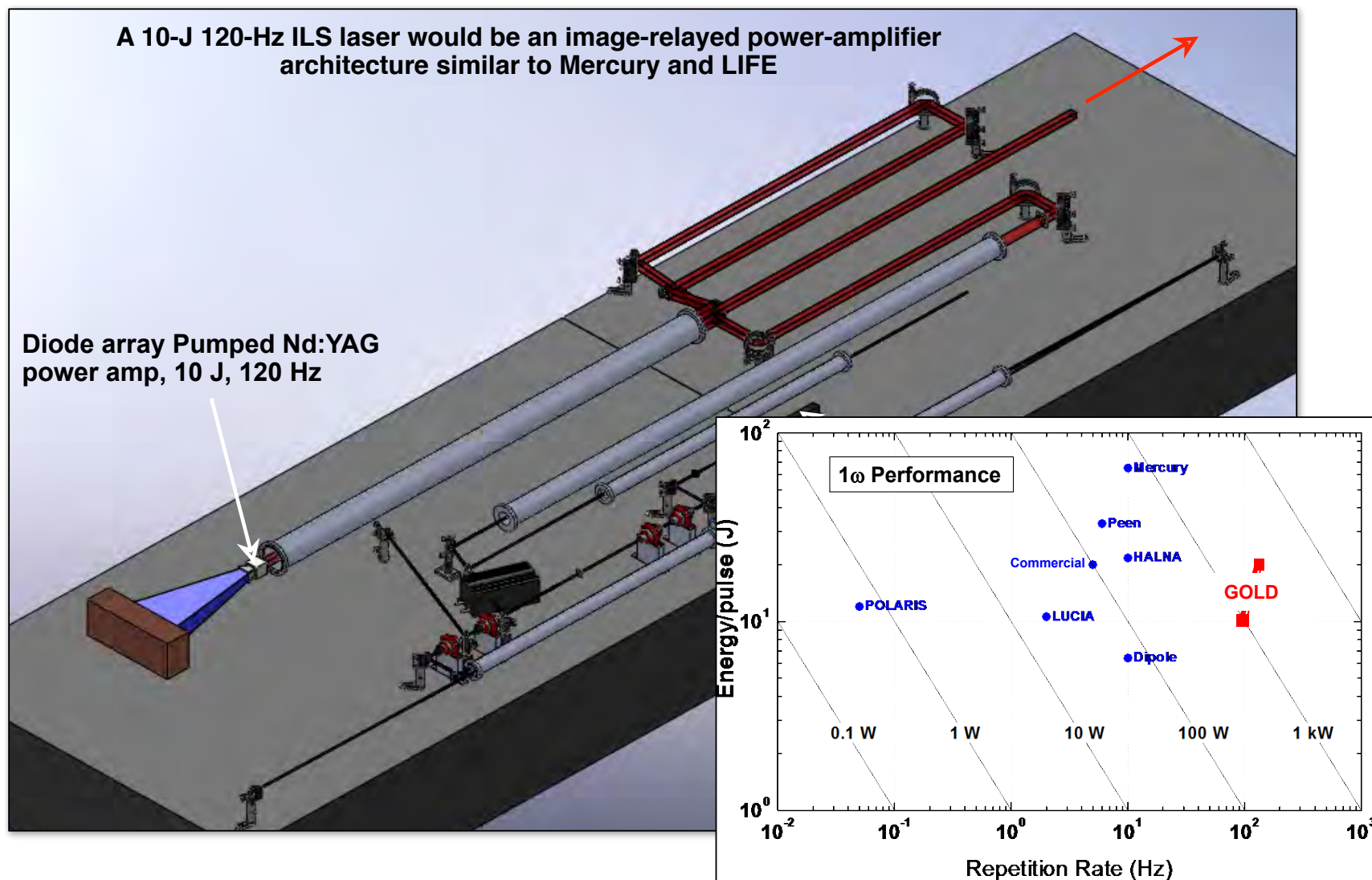
Rod Fluorescence showing gain uniformity



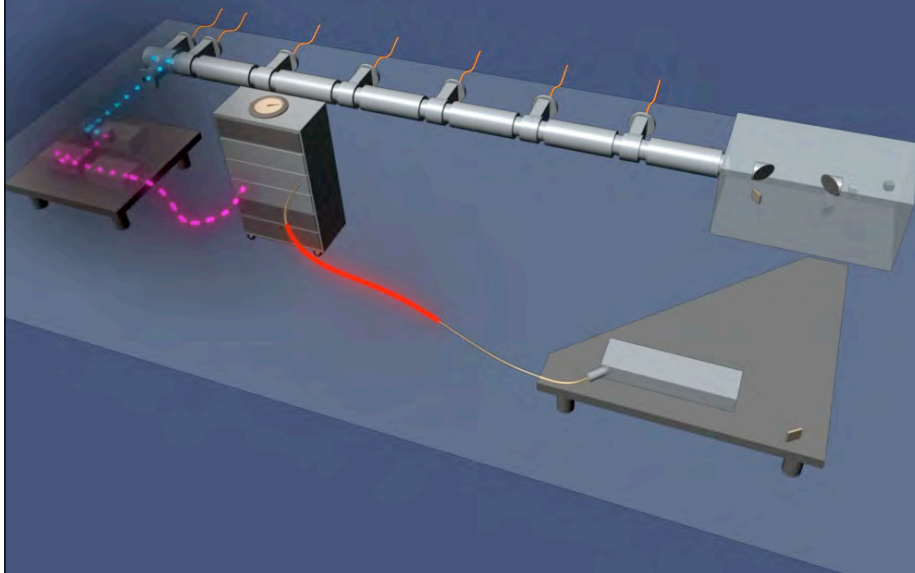
Our custom, diode-pumped solid state laser architecture is capable of $> 1\text{J}$ per pulse @ 120 Hz



LLNL is currently constructing a 120 Hz, 10 J, Nd:YAG laser for fusion optic lifetime studies



< 5 micro rad



**100% overlap
“Unity”**

**Conventional mirrors
NO interferometry required
Any color interaction
Intrinsically synchronized**

$\sim 10^{13}$ ph/s

